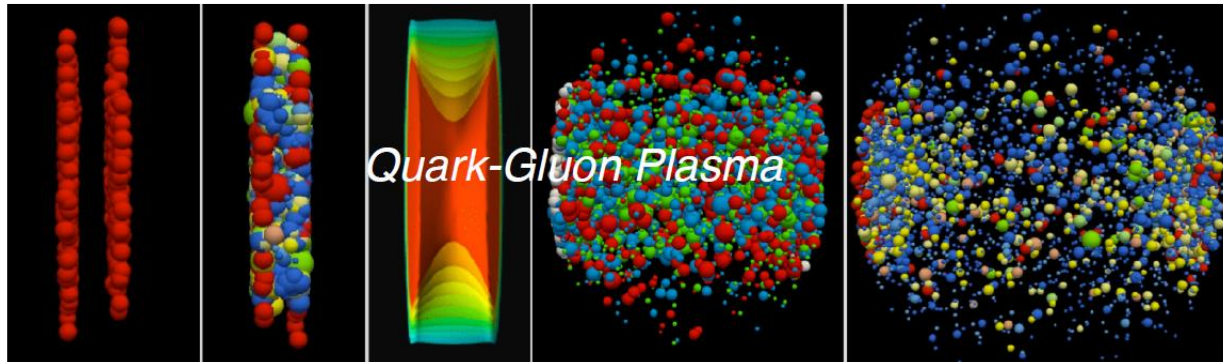


QCD Phase Transitions & Relativistic Heavy Ion Physics



- I. Why Heavy Ion Collisions — QCD matter***
- II. Quark-gluon plasma at RHIC and LHC***
- III. Rare particles (J/ψ , Υ , B_c , Ω_{ccc} , $cc\bar{c}\bar{c}$, $\mu^+\mu^-$)***

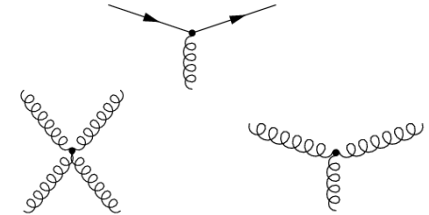
I. Why heavy ion collisions – QCD matter

QCD

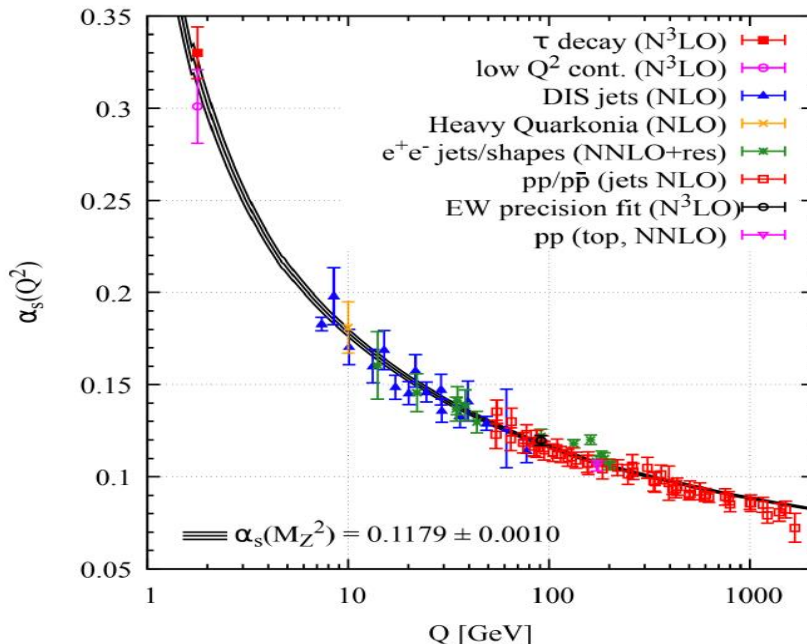
standard model $\left\{ \begin{array}{l} \text{Electromagnetic interaction} \\ \text{Weak interaction} \\ \text{Strong interaction, Quantum chromodynamics (QCD)} \end{array} \right.$

$$\mathcal{L}_{QCD} = \bar{\psi}_i (i\gamma_\mu D_{ij}^\mu - m\delta_{ij}) \psi_j - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu}$$

$$D_{ij}^\mu = \delta_{ij} \partial^\mu + i g A_\mu^a T_{ij}^a \quad G_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a - g f^{abc} A_\mu^b A_\nu^c$$



Non-Abelian interaction $\rightarrow$ QCD running coupling



Short range: weak coupling, asymptotic freedom, **pQCD**

2004 诺贝尔物理奖



David J. Gross



H. David Politzer



Frank Wilczek

Long range: strong coupling, color confinement, **non-perturbative properties !**

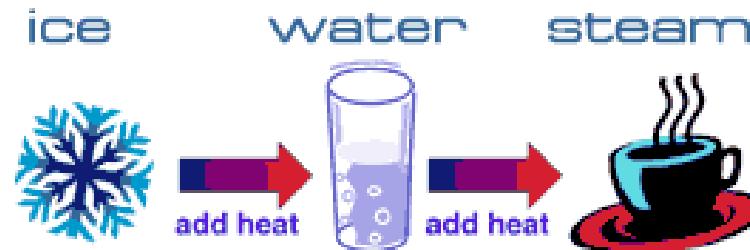
物质的状态

问题：看不见的夸克, 质量起源, 真空结构等非微扰效应。

方法：真空不知道，动动真空怎么样（例如量子力学中的**Casimir**效应）？

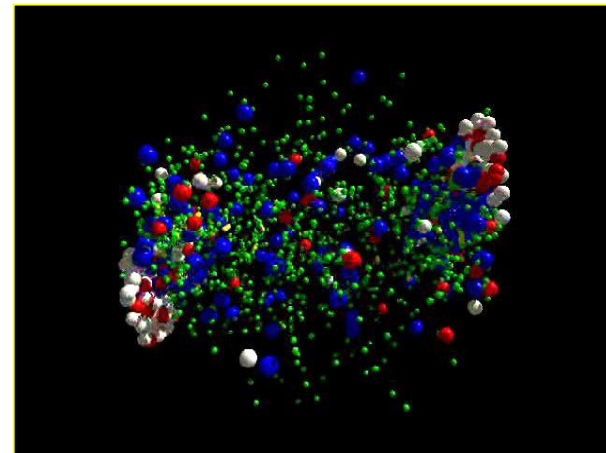
物质结构 → （相对于真空的）物质状态

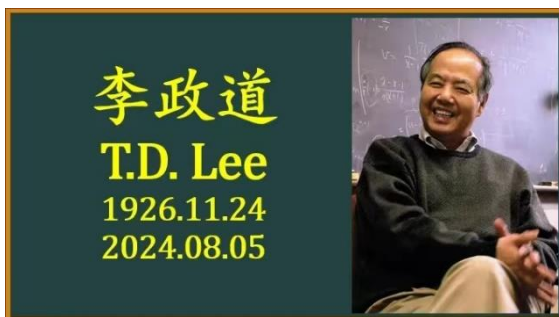
熟悉的物质三态：固态, 液态, 气态



第四态：原子弹爆炸后产生大范围的电磁等离子体

猜测：是否存在由夸克和胶子组成的**QCD**等离子体(**QGP**)？





李先生是相对论重离子碰撞物理的开创者！

统一研究物质的结构与状态

李政道【Rev. Mod. Phys. 47, 267(1975)】

In order to study the question of 'vacuum', we must turn to a different direction: we should investigate some 'bulk' phenomena by distributing high energy over a relatively large volume."

李政道【布鲁克海文报告, 2006】

*from 1897 to present: to comprehend the largest, we need only understand the smallest.
from present to 21st century: to know the smallest, we need also the largest!*

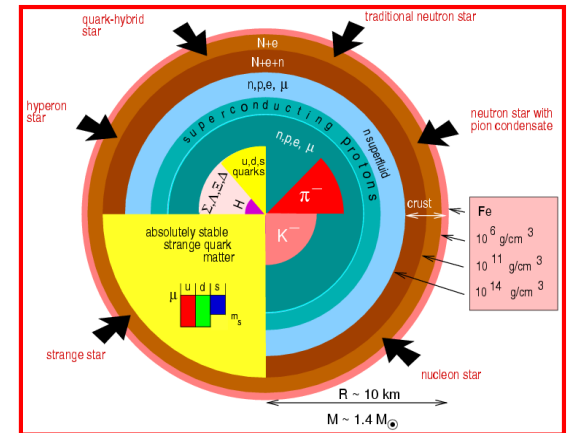
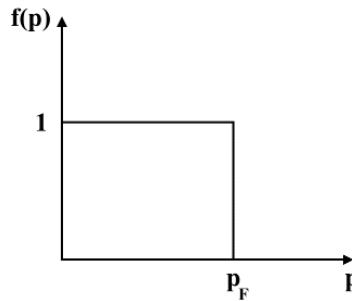
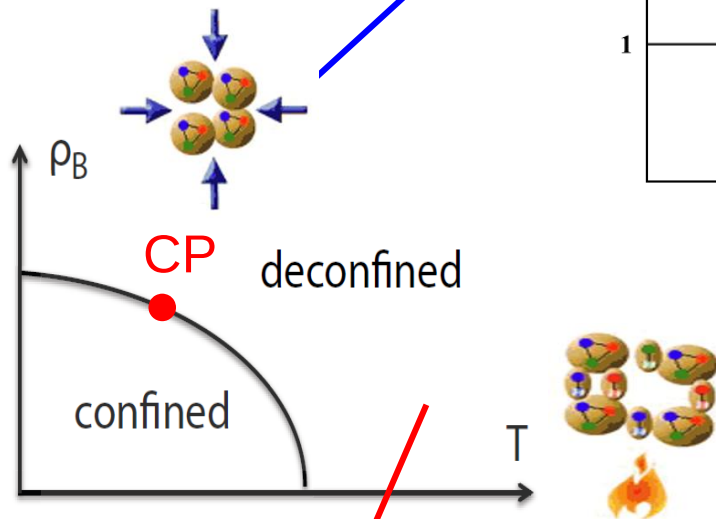
李政道【人民大会堂报告, 2012】

● 21世纪的物理发展
Holism 整体统一
微观的基本粒子和
宏观的真空态
需统一研究

Vacuum Excitation VS Matter Compression

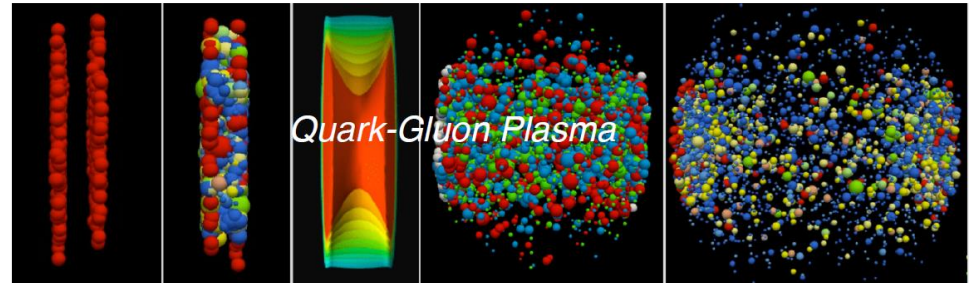
- ρ_B -induced QUANTUM Phase Transition
- Poor Theories

Matter compression !



Vacuum excitation !

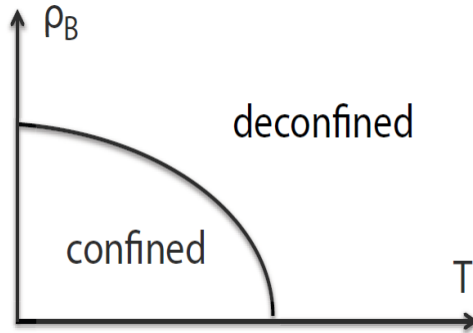
- T-induced CLASSICAL Phase Transition (Landau)
- Lattice QCD, (resumed)pQCD, FRG, DSE, Models



QCD Phase Diagrams

Deconfinement

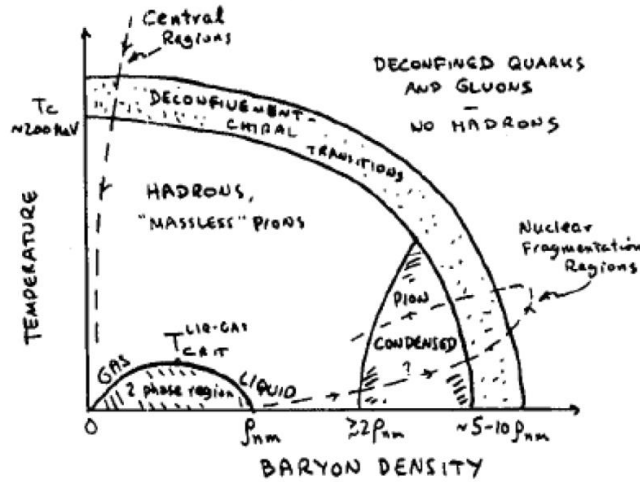
Cabbibo and Parisi, PLB59, 67(1975)



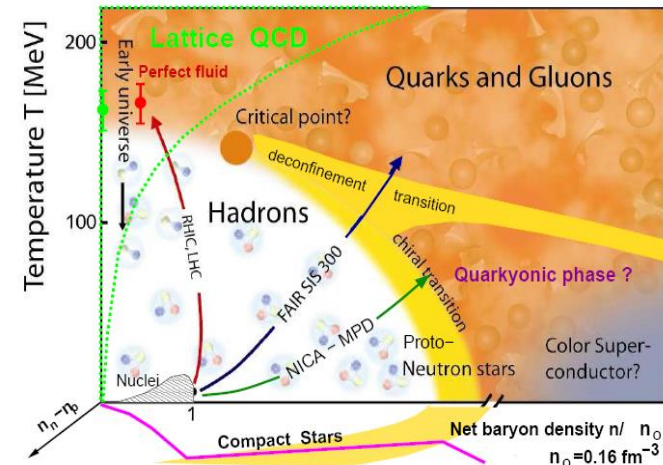
+ Chiral restoration

Baym, NSAC Long Rang Plan, 1983

PHASE DIAGRAM OF NUCLEAR MATTER



+ CSC, quarkyonic phase and critical point,



有哪些QCD相变?

相变是体系对称性质的改变

Deconfinement:

Chiral restoration:

Color superconductivity:

Pion superfluidity:

$Z(3)$ symmetry

$SU(3)$ chiral symmetry

$SU(3)$ color symmetry

$SU(2)$ isospin symmetry

研究量子物质的方法：量子统计力学

波函数的时间演化： $|\psi, t\rangle = \hat{U}(t, 0)|\psi, 0\rangle$

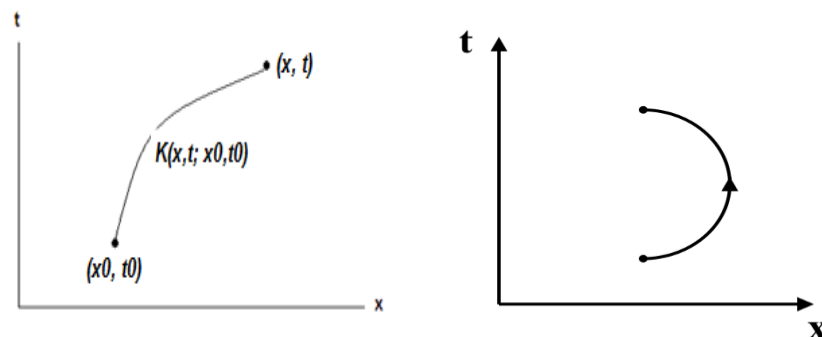
坐标表象： $\psi(x, t) = \langle x|\psi, t\rangle = \langle x|\hat{U}(t, 0)|\psi, 0\rangle$

$$= \int dx_0 \langle x|\hat{U}(t, 0)|x_0\rangle \langle x_0|\psi, 0\rangle = \int dx_0 K(x, t; x_0, 0) \psi(x_0, 0)$$

时空中波函数的传播子 $K(x, t; x_0, 0)$

空间封闭传播子 $K(x, t; x, 0)$

考虑所有的封闭传播子的集合：



$$\int dx K(x, t; x, 0) = \int dx \langle x|e^{-\frac{i}{\hbar}\hat{H}t}|x\rangle = \sum_E \int dx \langle x|E\rangle \langle E|e^{-\frac{i}{\hbar}\hat{H}t}|x\rangle$$

$$= \sum_E \int dx \langle E|x\rangle \langle x|E\rangle e^{-\frac{i}{\hbar}Et} = \sum_E e^{-\frac{i}{\hbar}Et} \xrightarrow{it/\hbar \rightarrow 1/T} \sum_E e^{-\frac{E}{T}} = Z$$

封闭传播子就是配分函数。用虚时动力学方法计算粒子的热力学。

极高温极高密: $pQCD$ Thermodynamics

Free QGP (封闭自由传播子):

$$\Omega_0(T, \mu) = -2N_c N_f \left[\frac{7\pi^2}{360} T^4 + \frac{\mu^2 T^2}{12} + \frac{\mu^4}{24\pi^2} \right] - 2(N_c^2 - 1) \frac{\pi^2}{90} T^4$$

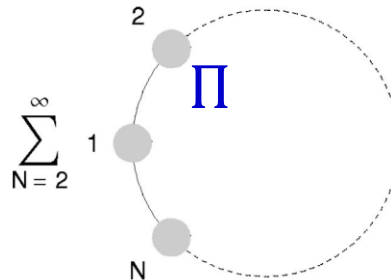
$\sim g$:

$$\Omega_1(T, \mu) = 0$$

$\sim g^2$:

$$\Omega_2^q(T, \mu) = \frac{N_f(N_c^2 - 1)}{8\pi} \alpha_s \left[\frac{5\pi^2}{18} T^4 + \mu^2 T^2 + \frac{\mu^4}{2\pi^2} \right]$$

Summation over ring diagrams:



HTL:

$$V_{HTL}(r) = -\tilde{\alpha}_s \left[m_D + \frac{e^{-m_D r}}{r} + iT\phi(m_D r) \right] + \mathcal{O}(g^4)$$

$$\phi(x) = 2 \int_0^\infty dz \frac{z}{(z^2 + 1)^2} \left(1 - \frac{\sin(xz)}{xz} \right)$$

袋模型估计相变温度

将QGP看成一个大口袋，袋内微扰真空，是弱耦合QGP，袋外物理真空。两种真空之差表现为压强作用在口袋上，平衡袋内的压强，形成稳定的口袋。

设袋外压强常数B。体系的压强为

$$\begin{aligned} P &= P_{QGP} - B \\ &= 2N_c N_f \left[\frac{7\pi^2}{360} T^4 + \frac{\mu^2 T^2}{12} + \frac{\mu^4}{24\pi^2} \right] + 2(N_c^2 - 1) \frac{\pi^2}{90} T^4 - B \end{aligned}$$

如果袋内为强子气体，只考虑 π 介子时的压强，

$$P = \frac{3\pi^2}{90} T^4$$

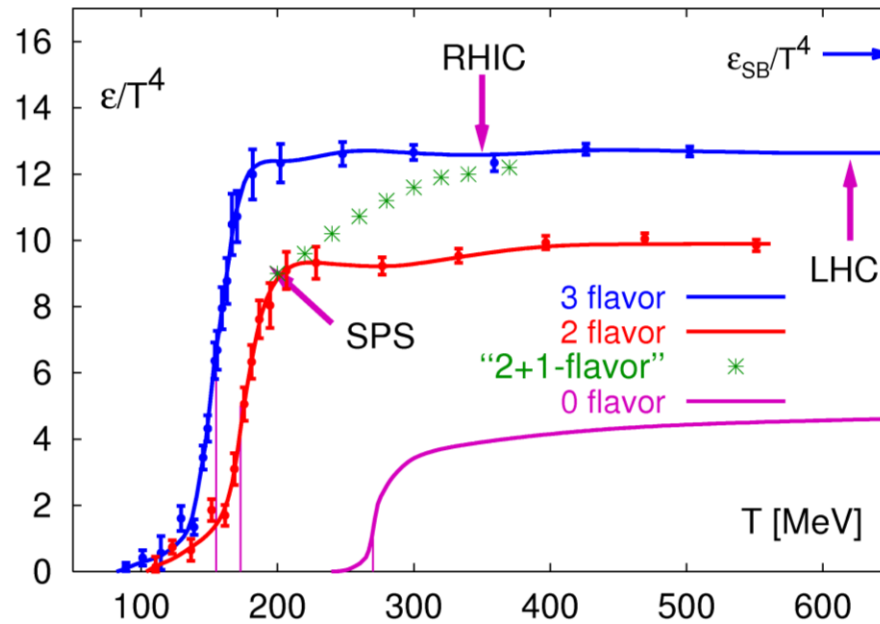
设体系在 T_c 时发生相变，由Gibbs相平衡条件，

$$2N_c N_f \left[\frac{7\pi^2}{360} T_c^4 + \frac{\mu^2 T_c^2}{12} + \frac{\mu^4}{24\pi^2} \right] + 2(N_c^2 - 1) \frac{\pi^2}{90} T_c^4 - B = \frac{3\pi^2}{90} T_c^4$$

取 $B^{1/4} = 200 \text{ MeV}$ ，有

$$T_c = 144 \text{ MeV}$$

有限温度低密: *Lattice QCD*

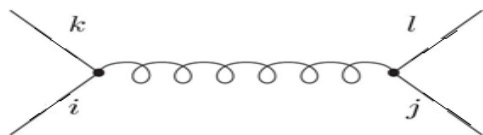


- *Not a phase transition but a crossover around $T_c \sim 160$ MeV*
- *Still interacting QGP at $T/T_c \sim 4$*

低温高密：色超导

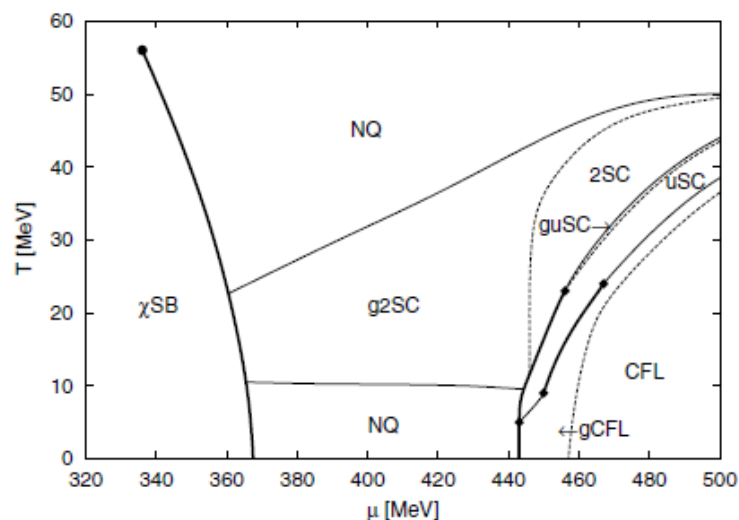
由BCS理论，在QED凝聚态中，两个电子通过交换光子是排斥相互作用，吸引相互作用是通过交换集体激发模式一声子来实现的 (低温超导)。Cooper对的形成自发破缺了电磁规范对称性，光子获得了质量。

Frank Wilczek(2004诺贝尔奖): 色超导研究的开创者之一



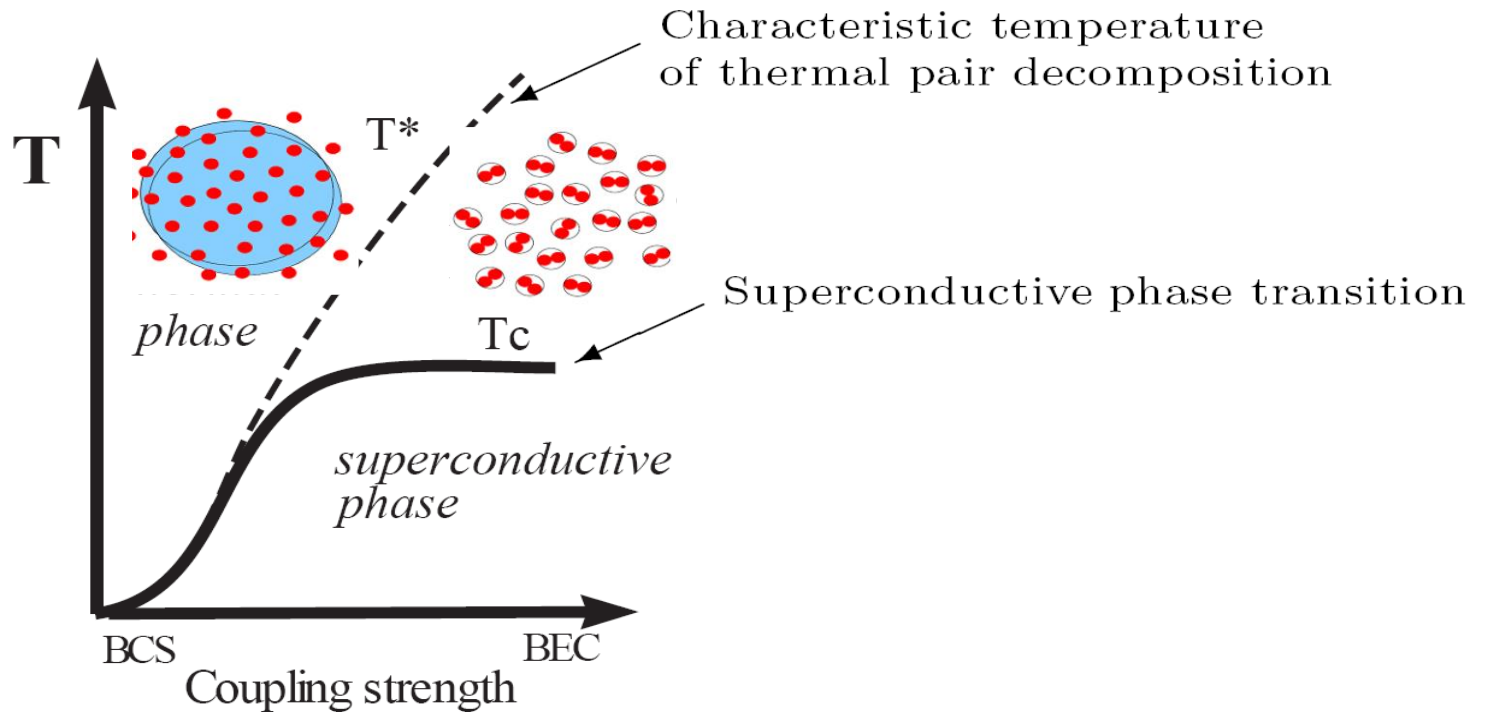
$$\begin{aligned} \sim (T_a)_{ki} (T_a)_{lj} = & -\frac{N_c + 1}{4N_c} (\delta_{jk} \delta_{il} - \delta_{ik} \delta_{jl}) \\ & + \frac{N_c - 1}{4N_c} (\delta_{jk} \delta_{il} + \delta_{ik} \delta_{jl}) \end{aligned}$$

第一项在交换初态或末态的两个夸克色指标时是反对称的，吸引相互作用。在单胶子交换的层次就使得两个夸克可以配对，**Diquark凝聚，形成色超导 (高温超导)**。由于色对称性自发破缺，胶子获得质量。



夸克的色味结构导致丰富的色超导相图

BCS-BEC Crossover



*in BCS, T_c is determined by thermal excitation of fermions,
in BEC, T_c is controlled by thermal excitation of collective Bosons.*

BCS-BEC crossover is probably a way to understand deconfinement in QCD !

Quantum Phase Transitions

Symmetry of $\mathcal{L}$
(may be explicitly broken like the term $m\bar{\psi}\psi$ in QCD)



$T = 0$: *spontaneously broken at classical level*
(mean field level), like chiral symmetry



$T \neq 0$: *classical phase transition,*
like chiral restoration



anomalously broken at quantum level
(loop level), like $U_A(1)$ symmetry

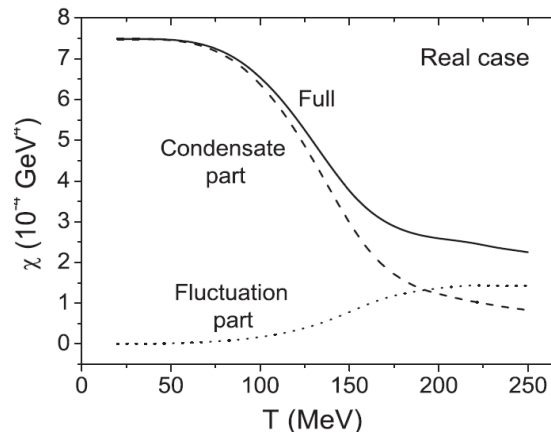
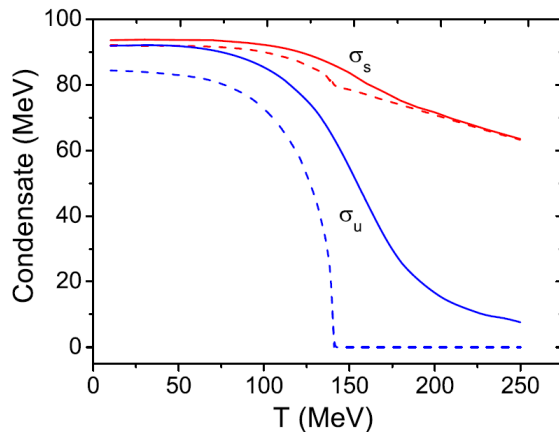


quantum phase transition,
like $U_A(1)$ restoration

• $U_A(1)$ symmetry

$$\chi = \int d^4x \langle T(Q(x)Q(0)) \rangle_{connected} \begin{cases} \neq 0 & U_A(1) \text{ breaking} \\ = 0 & U_A(1) \text{ symmetry} \end{cases}$$

$$Q(x) = \frac{g^2}{32\pi^2} F_{\mu\nu}^a(x) \tilde{F}_a^{\mu\nu}(x)$$



• *Chiral symmetry is restored at $T_c \sim 150$ MeV, but $U_A(1)$ symmetry is still broken above T_c !*

Lattice QCD:

HotQCD, PRD86, 094503(2012)

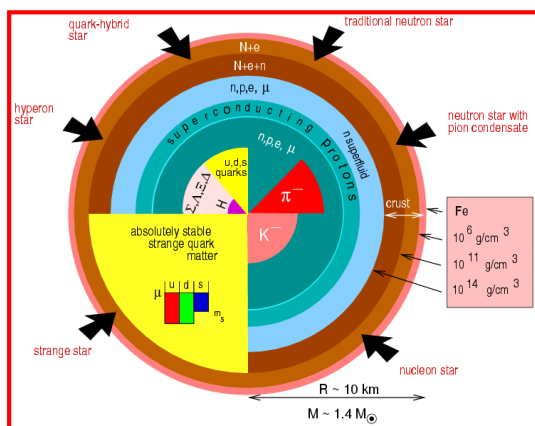
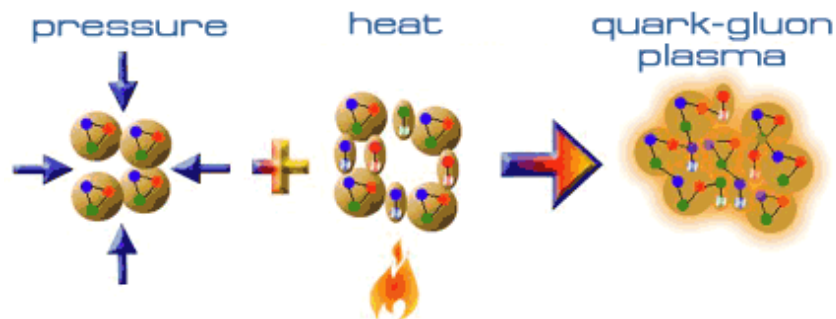
Effective models:

Jiang, Xia and PZ, PRD93, 074006(2016)

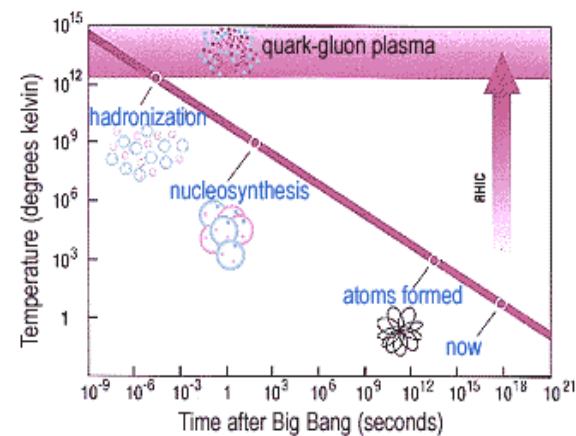
Li, Gui, Zhuang, CPC, 2023

如何实现QGP

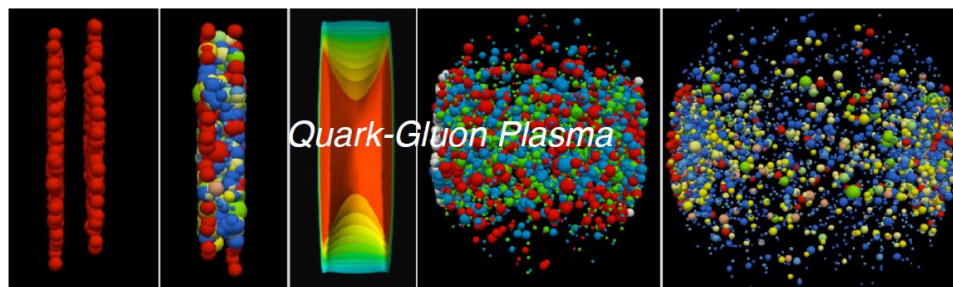
产生QGP的条件:



致密星体



早期宇宙



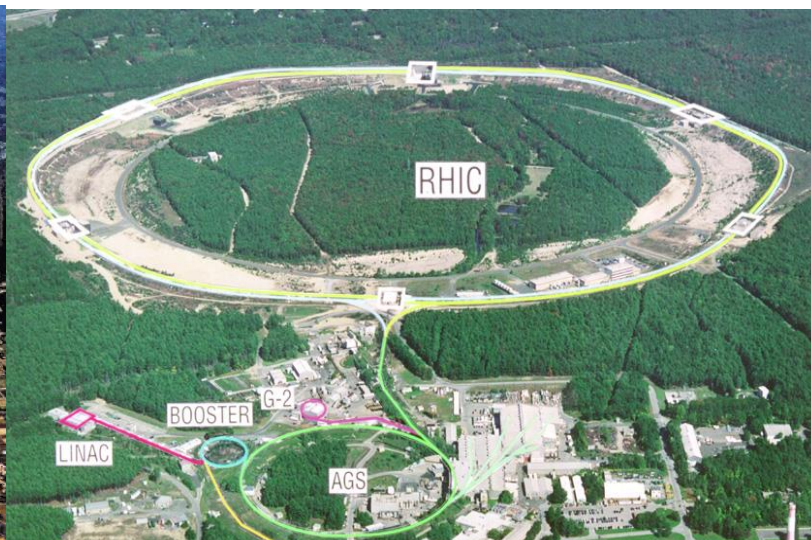
相对论重离子碰撞是在实验室产生QGP的唯一可能手段!

相对论重离子碰撞是在实验室实现QGP的唯一可能方式

LHC, Pb+Pb@ $\sqrt{s} = 5500$ GeV, *ALICE, ATLAS, CMS, LHCb*



RHIC, Au+Au@ $\sqrt{s} = 200$ GeV, *STAR, PHENIX*



FAIR, U+U@ $E_{lab} = 40A$ GeV, *CBM*



CSR, U+U@ $E_{lab} = 0.6A$ GeV, *CEE*

常用物理量

核几何

碰撞参数 $\vec{b}$: 两个碰撞核中心的相对位置, 矢量。 $\vec{b} = 0$, 中心碰撞。

参与碰撞核子数 N_{part} : 软过程数 $\sim N_{part}$

参与碰撞的核子对数 N_{coll} : 硬过程数 $\sim N_{coll}$

纵向快度 (沿碰撞方向) $y = \frac{1}{2} \ln \frac{E+p_z}{E-p_z}$, $\begin{cases} E = m_{\perp} \cosh(y) \\ p_z = m_{\perp} \sinh(y) \end{cases}$

优点: Lorentz变换时, 快度相加 $y' = y + Y$

E_{lab} 与 $\sqrt{s}$ 的转换

$$\sqrt{s} = \sqrt{(E_1 + E_2)^2 - (\vec{p}_1 + \vec{p}_2)^2} = \sqrt{(E_{lab} + m_N)^2 - \vec{p}_1^2} = \sqrt{2m_N(m_M + E_{lab})}$$

正常核物质的能量密度

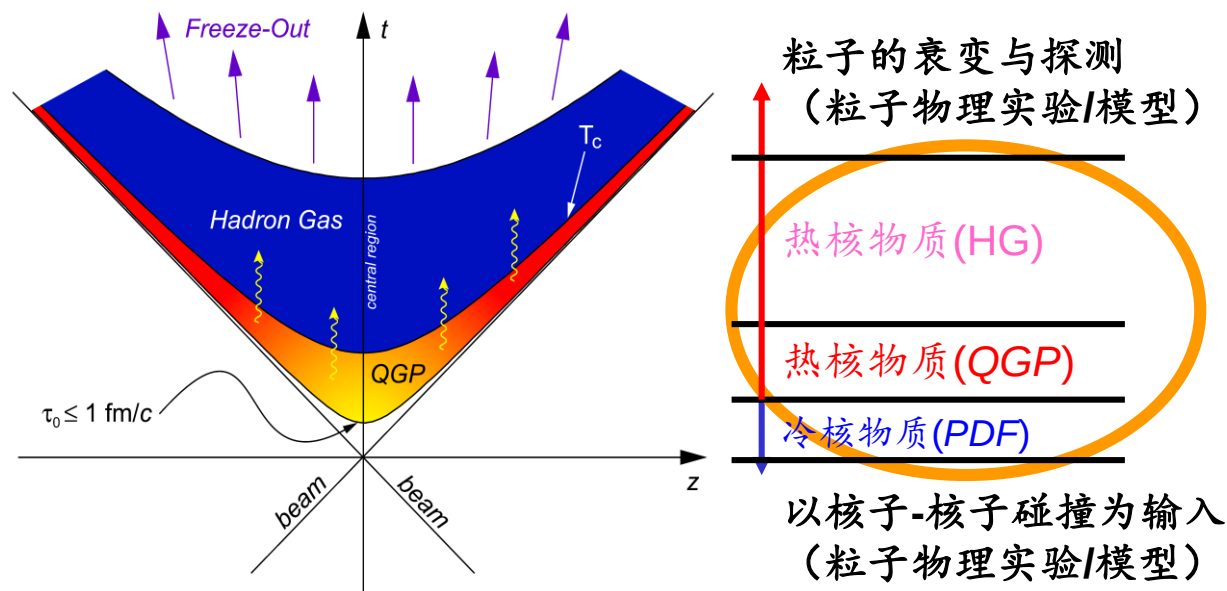
核内核子数密度 $n_0 = 0.17/fm^3$, 核内能量密度 $\varepsilon_0 = m_N n_0 \sim 0.17 GeV/fm^3$

产生QGP的最低能量密度

$$\text{即核子内的能量密度 } \varepsilon_{min} = \frac{m_N}{\frac{4}{3}\pi r_N^3} \xrightarrow{r_N=0.8 fm} 0.5 GeV/fm^3$$

QGP的信号

QGP只是重离子碰撞的中间态，即使在RHIC和LHC产生了QGP，也不能在末态直接观测到，只能通过携带QGP信号的末态粒子来判断。



希望看到什么？

1) 与 pp 的不同 (*off-diagonal interactions*), 有新的物理

$$\text{核修正因子: } R_{AA} = \frac{\sigma_{AA}}{N_{coll}\sigma_{pp}}$$

2) 与强子气体的差别, 夸克层次的新物理, QGP的产生!

热密物质的时空演化—流体力学

Landau: 用流体力学描述高能碰撞后至衰变前体系的时空演化。

忽略粘滞效应的理想流体力学方程

$$\begin{cases} \partial_\mu T^{\mu\nu} = 0 & (\text{能动量守恒}) \\ \partial_\mu n^\mu = 0 & (\text{重子数守恒}) \end{cases}$$

$$T^{\mu\nu} = -Pg^{\mu\nu} + (\varepsilon + P)u^\mu u^\nu, \quad n^\mu = n_B u^\mu, \quad u_\mu u^\mu = 1$$

+ *Equation of state* $\varepsilon(P)$ (*QGP or Hadron gas from QCD thermodynamics*)

$\rightarrow u_\mu(x), T(x), \mu_B(x), P(x)$

一维膨胀 $\rightarrow$ 快度中心区的能量密度(*Bjorken*估计)

$$\begin{cases} \varepsilon(\tau) = \varepsilon(\tau_0) \left(\frac{\tau_0}{\tau}\right)^{1+c_s^2} \\ n(\tau) = n(\tau_0) \frac{\tau_0}{\tau} \\ T(\tau) = T(\tau_0) \left(\frac{\tau_0}{\tau}\right)^{c_s^2} \end{cases}$$

Probes of QCD Phases

