



# Quantum-Well-Controlled High-Order Anisotropic Magnetoresistance in Epitaxial Fe(001) Thin Films



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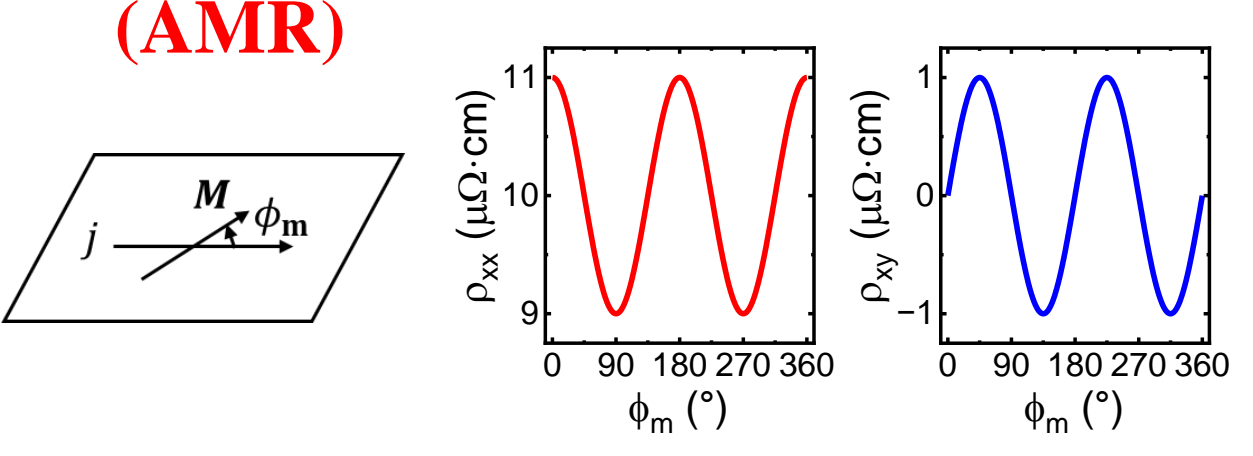
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## Introduction

### Anisotropic magnetoresistance (AMR)



$$\rho_{xx}(\phi) = \rho_0 + \Delta\rho_2 \cos 2\phi$$

$$\rho_{xy}(\phi) = \Delta\rho_2 \sin 2\phi \leftarrow \text{Planar Hall effect (PHE)}$$

J. M. D. Coey, Magnetism and Magnetic Materials

### Phenomenological model of AMR based on crystal symmetry

W. Döring, Ann. Phys. 424, 259 (1938) From Ohm's law  $E_i = \rho_{ij}J_j$ , expansion of  $\rho_{ij}$  on  $\mathbf{m}$ :

$$\rho_{ij}(\mathbf{m}) = a_{ij} + a_{klj}m_k m_l + a_{klmij}m_k m_l m_m m_n + \dots$$

Onsager relation:  $\rho_{ij}(\mathbf{m}) = \rho_{ji}(-\mathbf{m})$   
Cubic crystalline symmetry

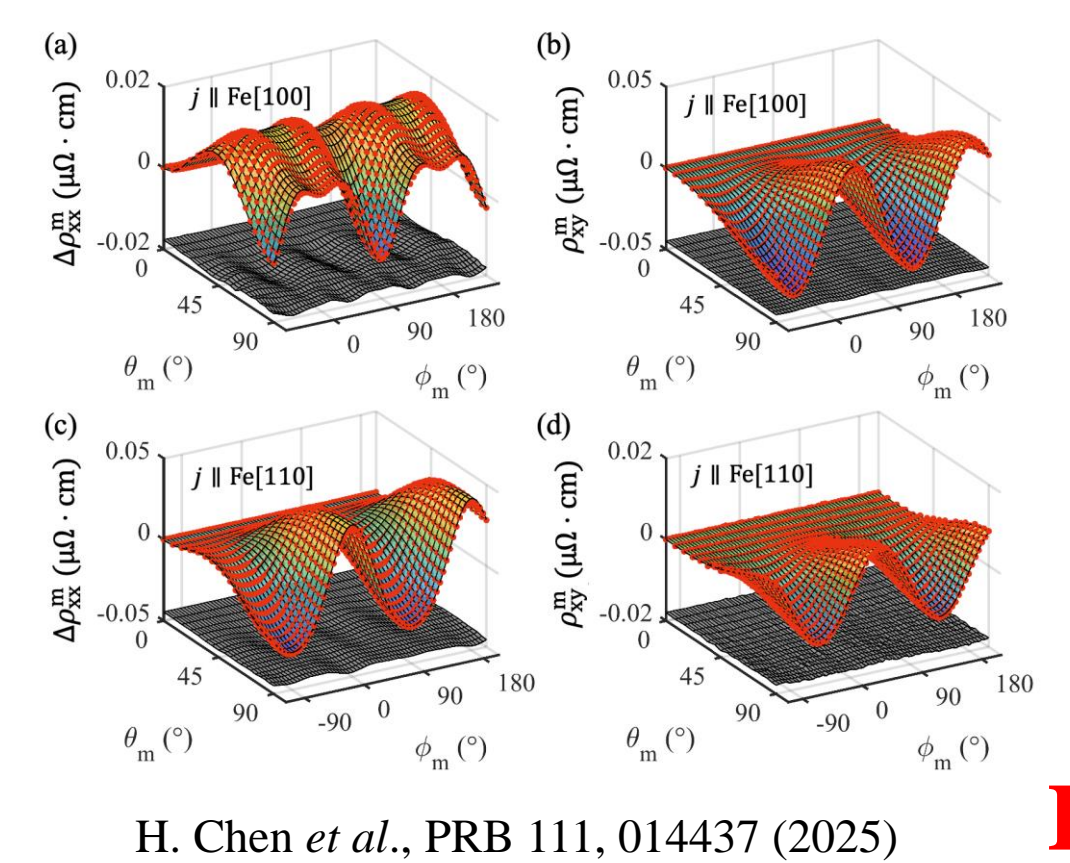
$$\rho_{xx}^{\parallel[100]}(\phi) = \rho_0 + \Delta\rho_2 \cos(2\phi) + \Delta\rho_4 \cos(4\phi)$$

$$\rho_{xx}^{\parallel[110]}(\phi) = \rho_0 + \Delta\rho_2' \cos(2\phi) - \Delta\rho_4 \cos(4\phi)$$

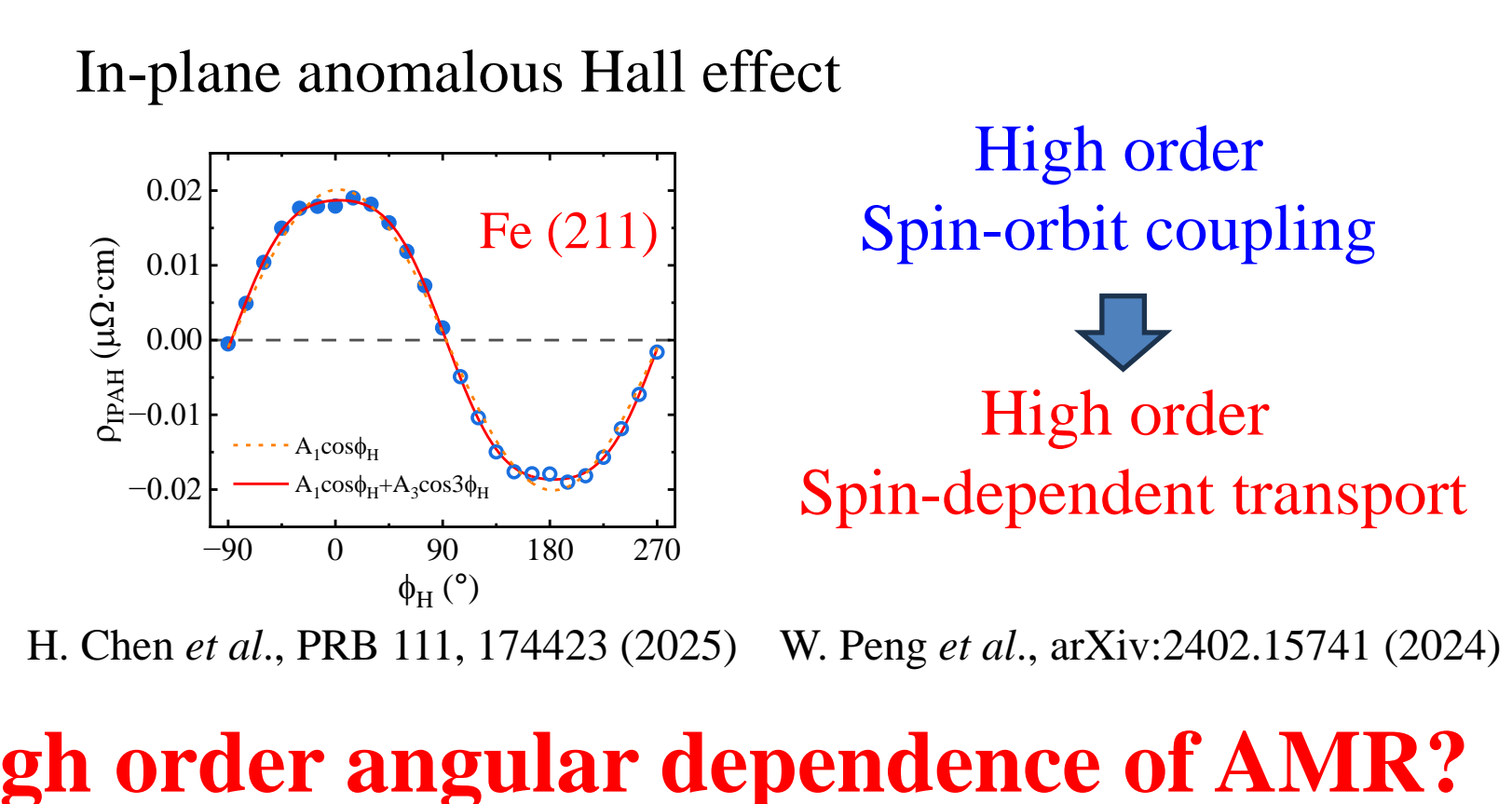
$$\rho_{xy}^{\parallel[100]}(\phi) = \Delta\rho_2 \sin(2\phi)$$

$$\rho_{xy}^{\parallel[110]}(\phi) = \Delta\rho_2' \sin(2\phi)$$

### Experimental verification



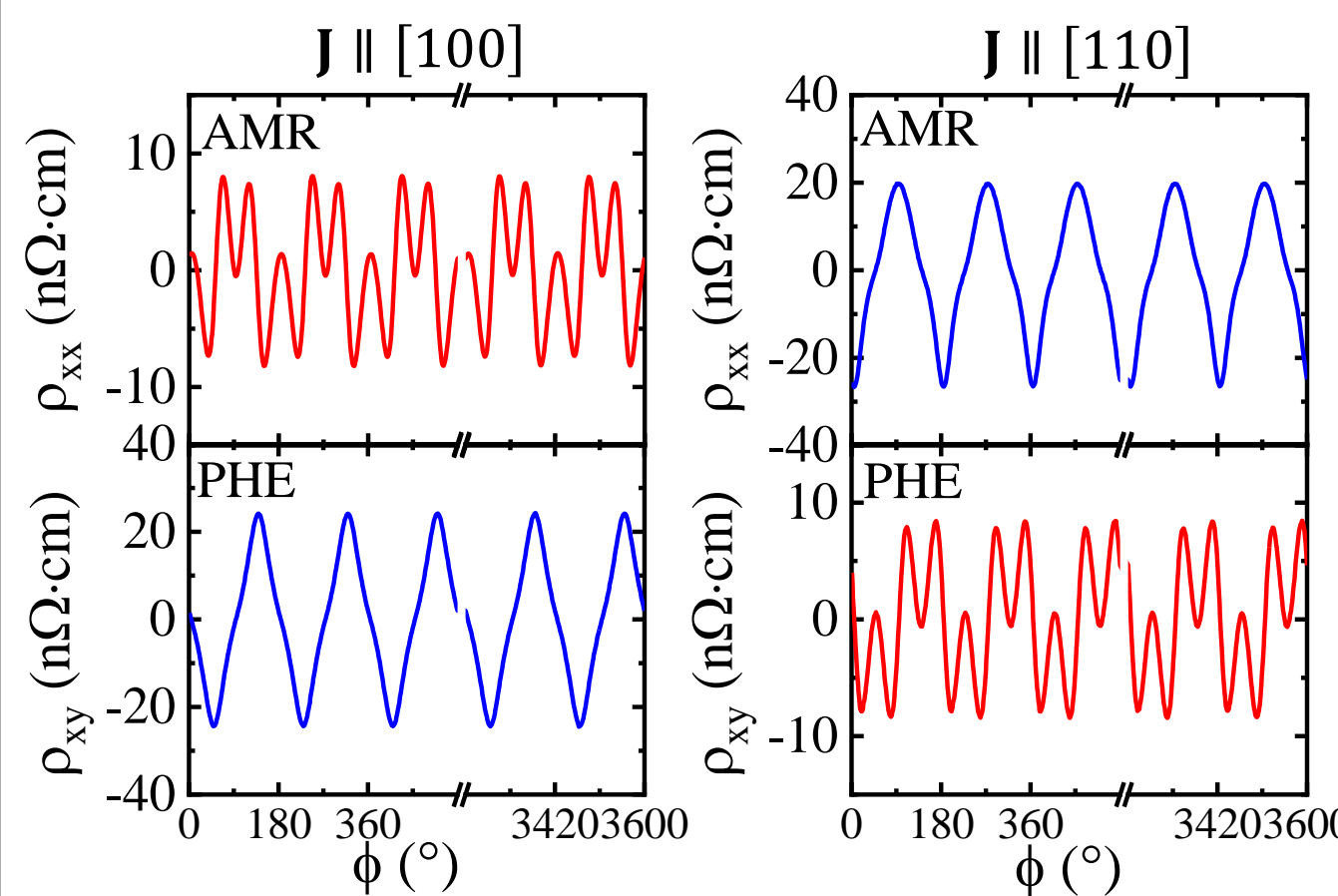
### High order galvanomagnetic effects



### High order angular dependence of AMR?

## High order angular-dependent AMR

### Reciprocal relation



High-order expansion of AMR and PHE: (Considering  $C_4$  symmetry, truncated at 6<sup>th</sup> order)

$$\rho_{xx}^{\parallel[100]}(\phi) = \Delta\rho_2 \cos(2\phi) + \Delta\rho_4 \cos(4\phi) + \Delta\rho_6 \cos(6\phi)$$

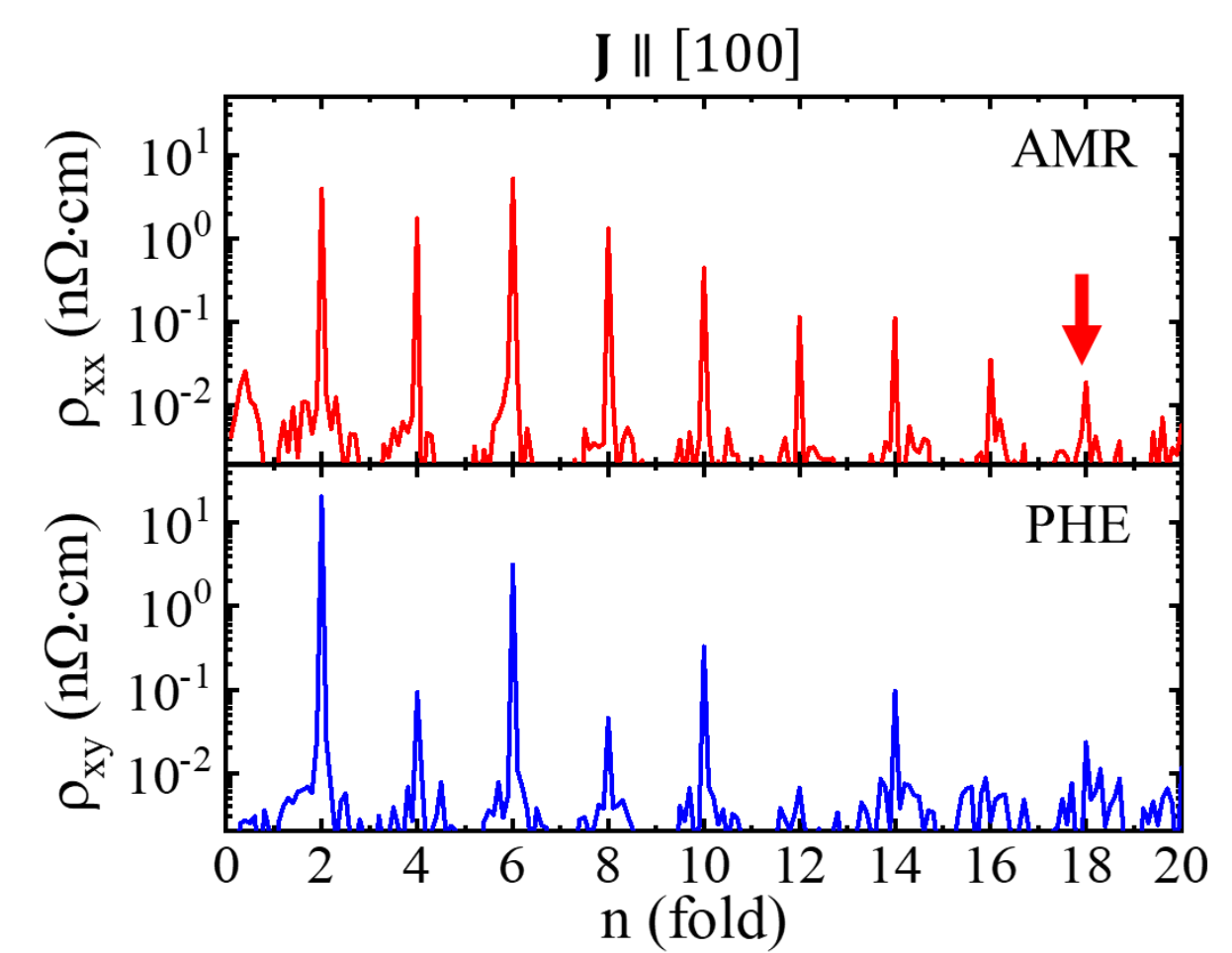
$$\rho_{xy}^{\parallel[100]}(\phi) = \Delta\rho_2' \sin(2\phi) - \Delta\rho_6' \sin(6\phi)$$

$$\rho_{xx}^{\parallel[110]}(\phi) = \Delta\rho_2' \cos(2\phi) - \Delta\rho_4 \cos(4\phi) + \Delta\rho_6' \cos(6\phi)$$

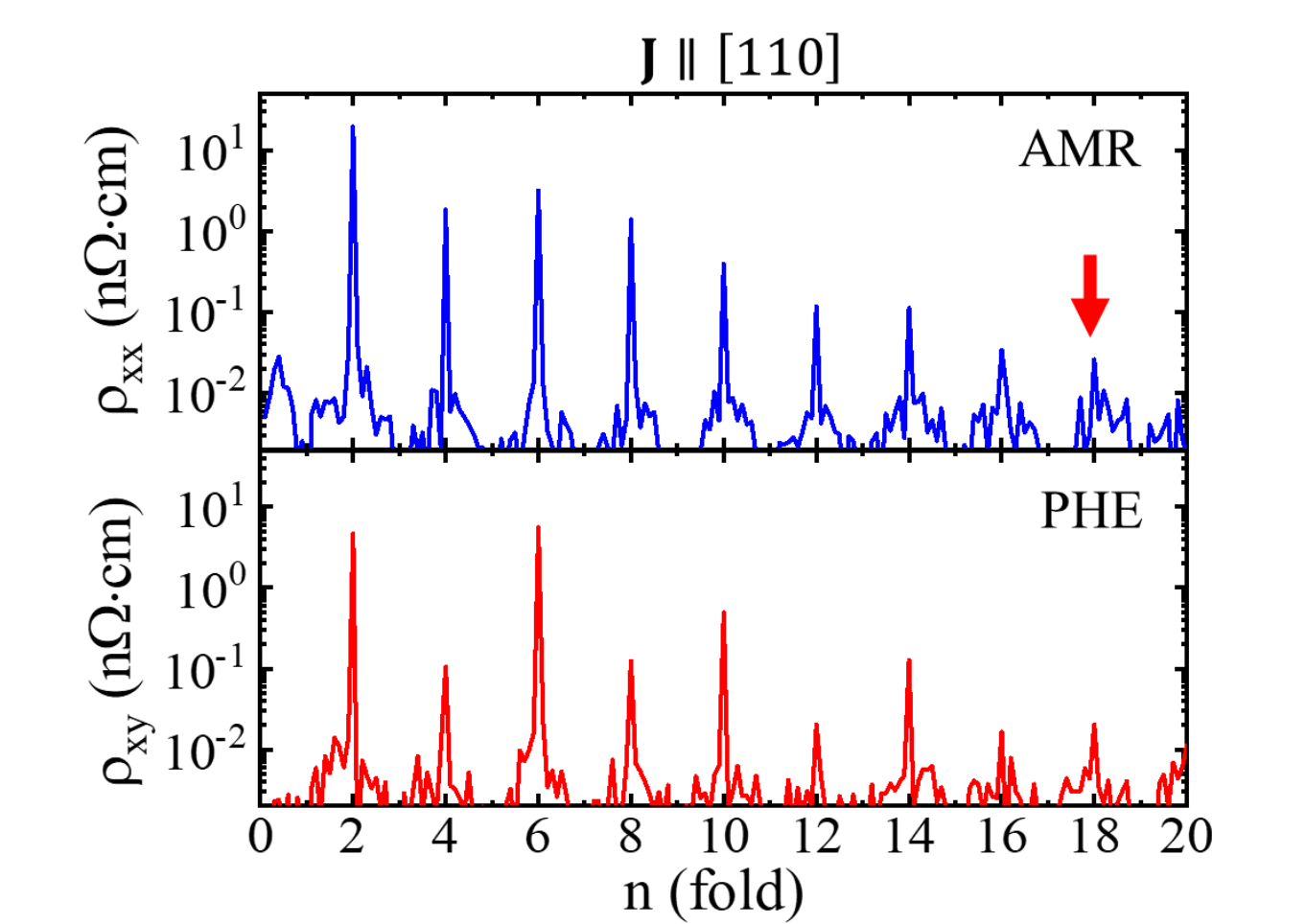
$$\rho_{xy}^{\parallel[110]}(\phi) = \Delta\rho_2 \sin(2\phi) - \Delta\rho_6 \sin(6\phi)$$

Consistent with experimental results

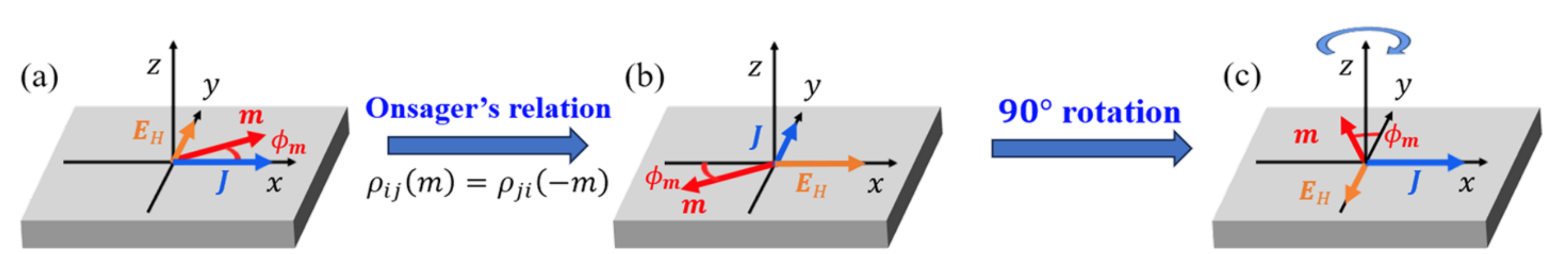
### Fast Fourier transformation



### Up to 18-fold AMR



### Onsager relation prohibits 4-fold PHE with perfect $C_4$ symmetry



Constraint on PHE:

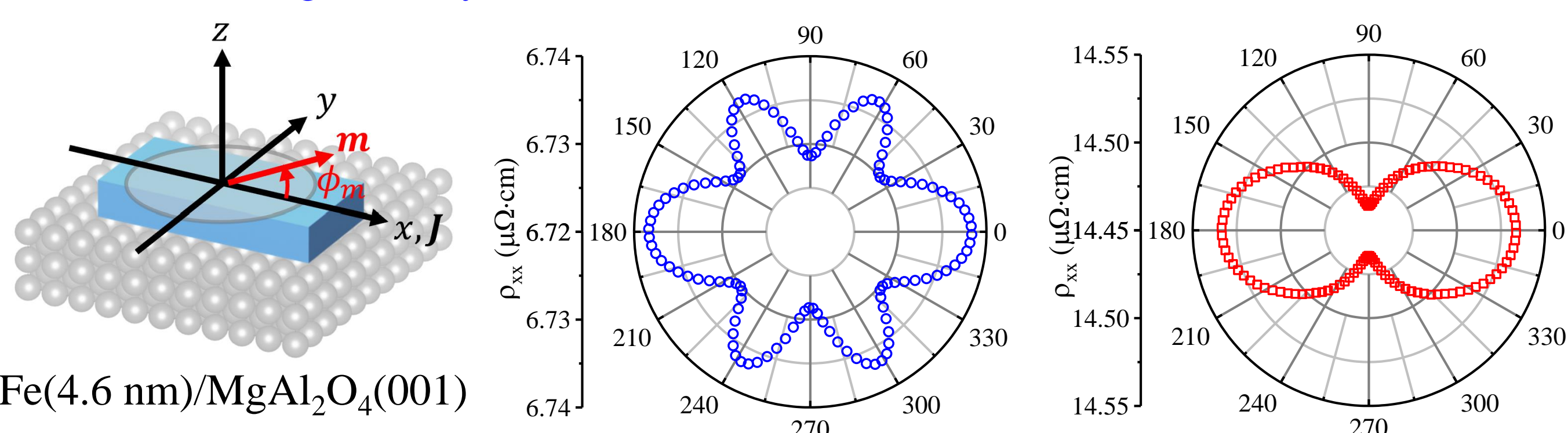
$$\rho_{xy}(\phi) = -\rho_{xy}(\phi + \pi/2)$$

$$\cos 2(\phi + \pi/2) = -\cos 2\phi \quad \checkmark$$

$$\cos 4(\phi + \pi/2) = \cos 4\phi \quad \times$$

## Six-fold angular-dependent AMR

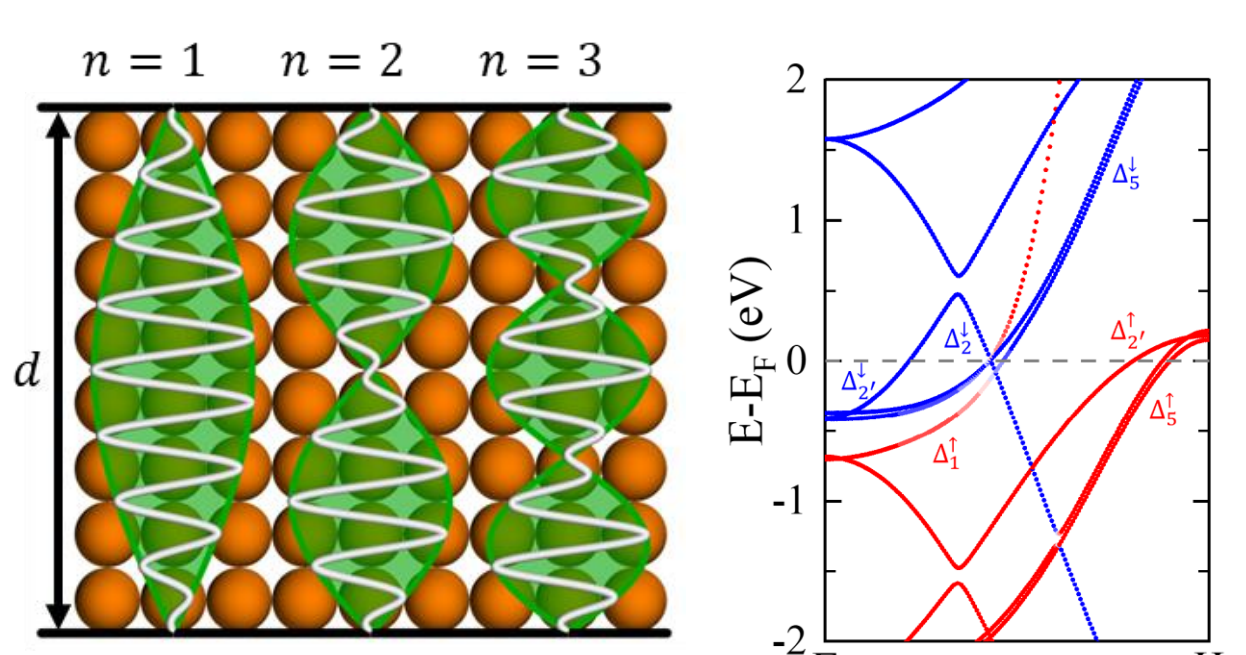
### Measurement geometry:



Six-fold angular-dependent AMR is observed at low temperature

## Quantum-well-induced oscillatory AMR in thin films

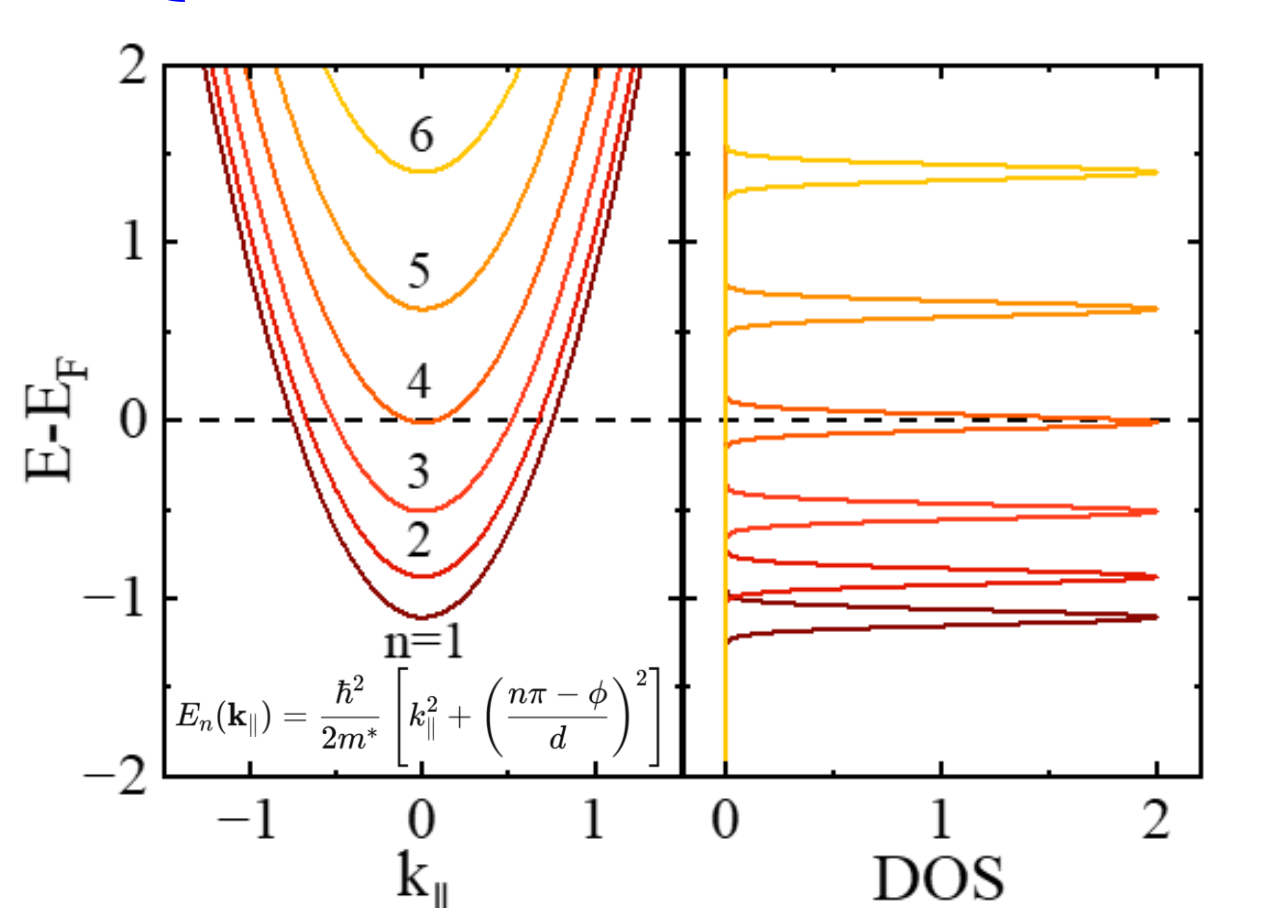
### Formation of quantum well states



Confinement along  $d \rightarrow$  Discrete states

$$\text{QWS condition: } 2k_F d + \varphi = 2n\pi$$

### QWS-induced AMR oscillations



Increase  $d \rightarrow$  QWS subbands cross  $E_F$

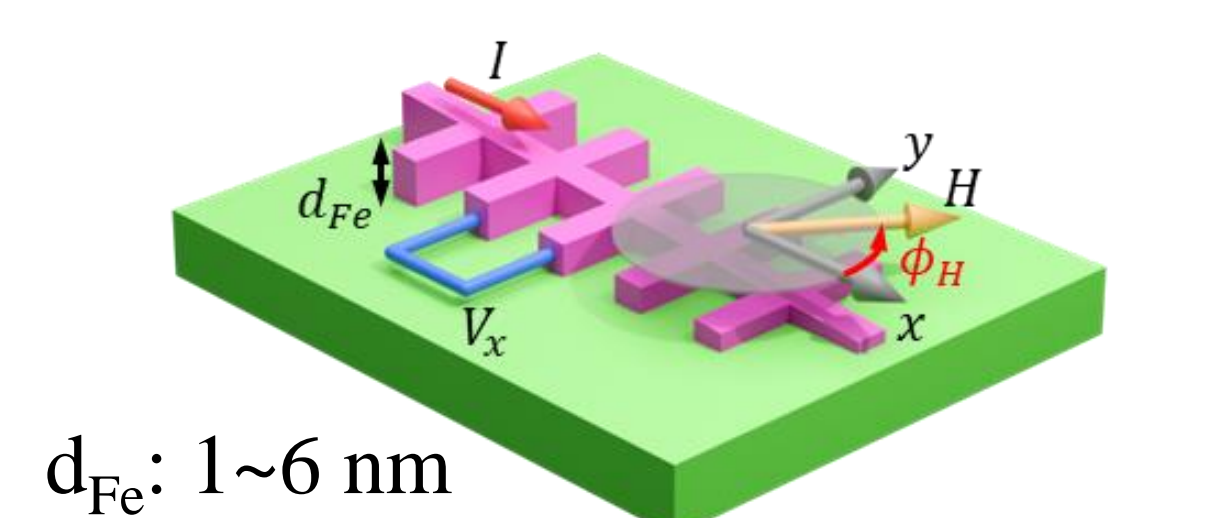
DOS at  $E_F$  oscillates with  $d$

AMR exhibits corresponding oscillation

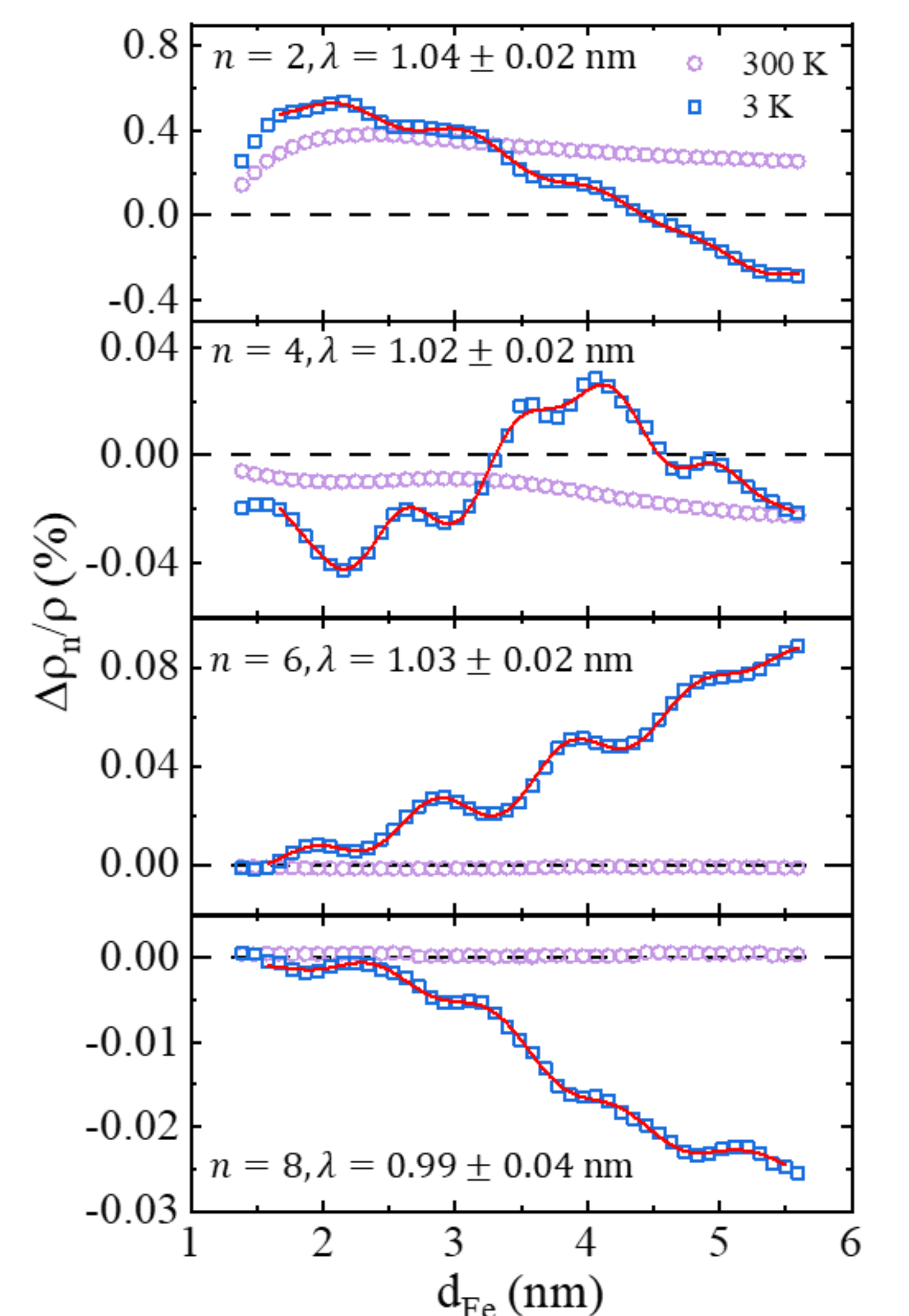
$$\text{Period: } \lambda = \frac{\pi}{k_F} \quad \text{or} \quad \frac{\pi}{k_{BZ} - k_F}$$

$$\text{For } \Delta_2^{\uparrow}, \text{ band: } \lambda \sim 1.0 \text{ nm}$$

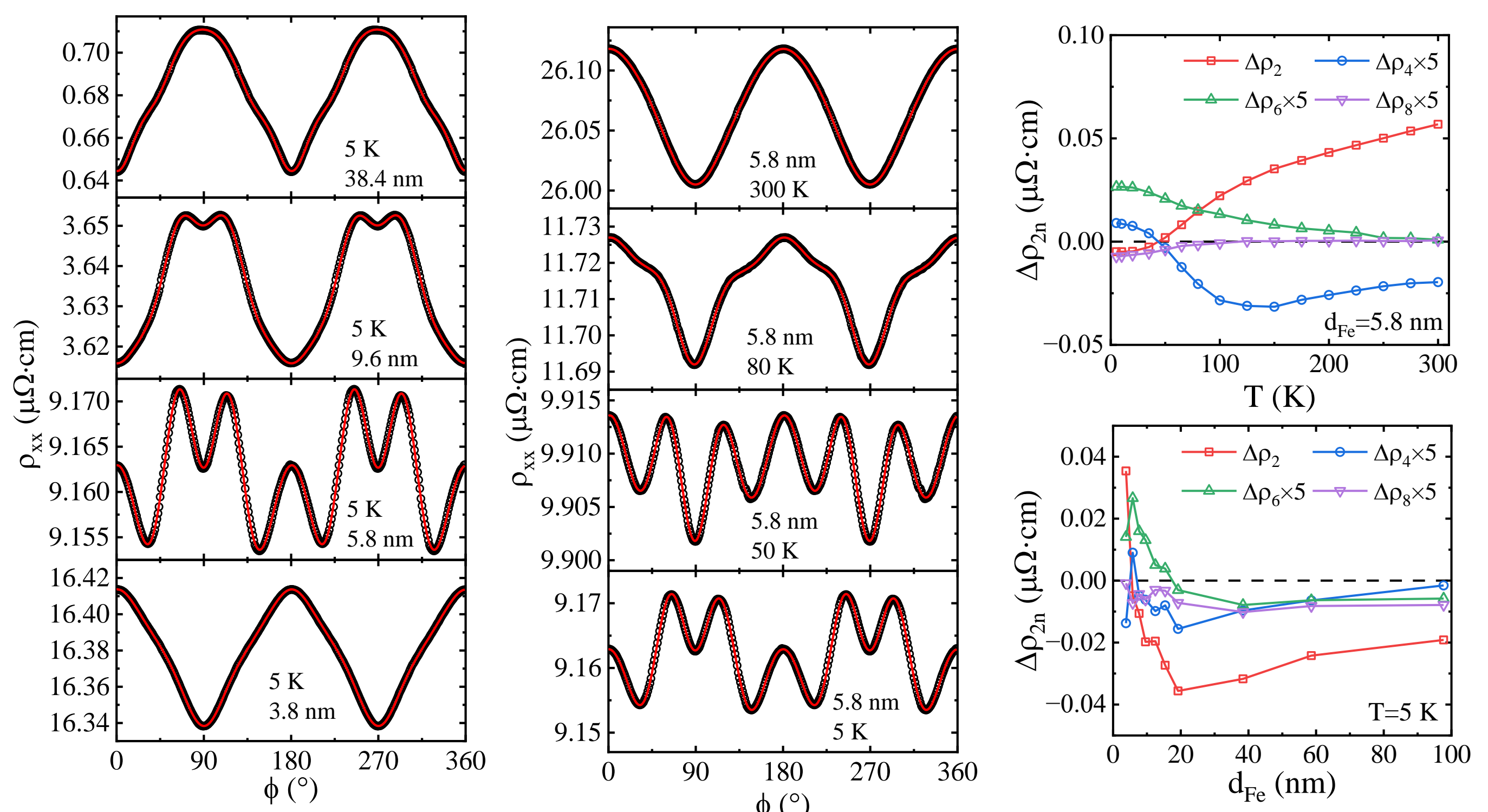
### Measurement in a wedge sample



$d_{Fe}$ : 1~6 nm

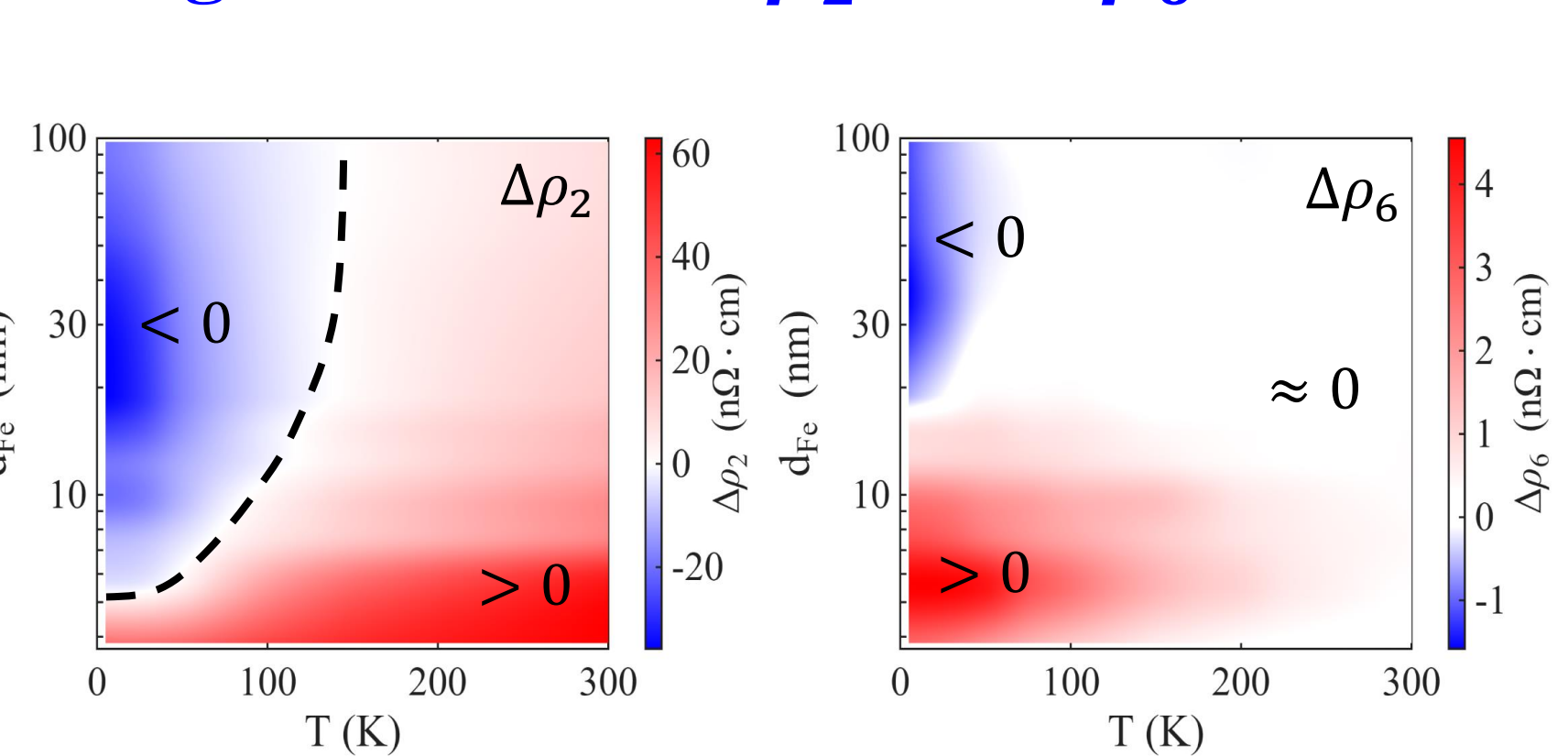


## Temperature- and thickness-dependent AMR



Fitting function:  $\rho_{xx}(\phi) = \rho_0 + \Delta\rho_2 \cos 2\phi + \Delta\rho_4 \cos 4\phi + \Delta\rho_6 \cos 6\phi + \dots$

### Sign reversal of $\Delta\rho_2$ and $\Delta\rho_6$



### Possible explanation

$$\sigma_{xx} \sim \sum_{nk} v_{x,nk}^2 \tau_{nk} = \bar{\tau}_0 \sum_{nk} v_{x,nk}^2$$

$$\rho_{xx} \sim \frac{1}{\bar{\tau}_0} \cdot \frac{1}{\sum_{nk} v_{x,nk}^2} = f^r(\phi) \cdot g^v(\phi)$$

Relaxation time contribution:  
 $f^r(\phi) = c_0^r + c_2^r \cos 2\phi + c_4^r \cos 4\phi$

Fermi velocity contribution:  
 $g^v(\phi) = c_0^v + c_2^v \cos 2\phi + c_4^v \cos 4\phi$

$$\Delta\rho_6 = 1/2 \cdot (c_4^v c_2^r + c_4^r c_2^v) \cos 6\phi$$

## Summary

- 6-fold angular AMR in nominally 4-fold Fe(001) films
- Sign reversal of 2-fold and 6-fold AMR with thickness
- Reciprocal relation between AMR and PHE
- AMR modulation by  $\Delta_2^{\uparrow}$ -band quantum well states

H. Chen *et al.*, Phys. Rev. Lett. 136, 086704 (2026)