

# Reflected Multientropy and Its Holographic Dual

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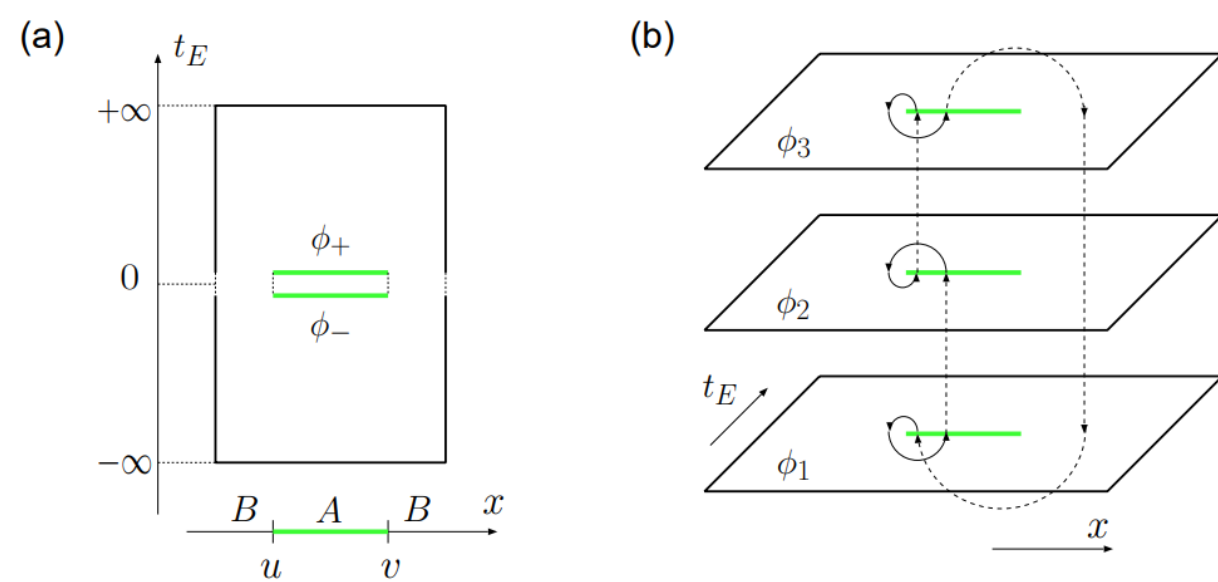


## 1. Rényi entropy and holographic dual

For a bipartite pure state, entanglement entropy is obtained from the Rényi entropy

$$S_n(A) = \frac{1}{1-n} \log \text{Tr} \rho_A^n, \quad S(A) = \lim_{n \rightarrow 1} S_n(A).$$

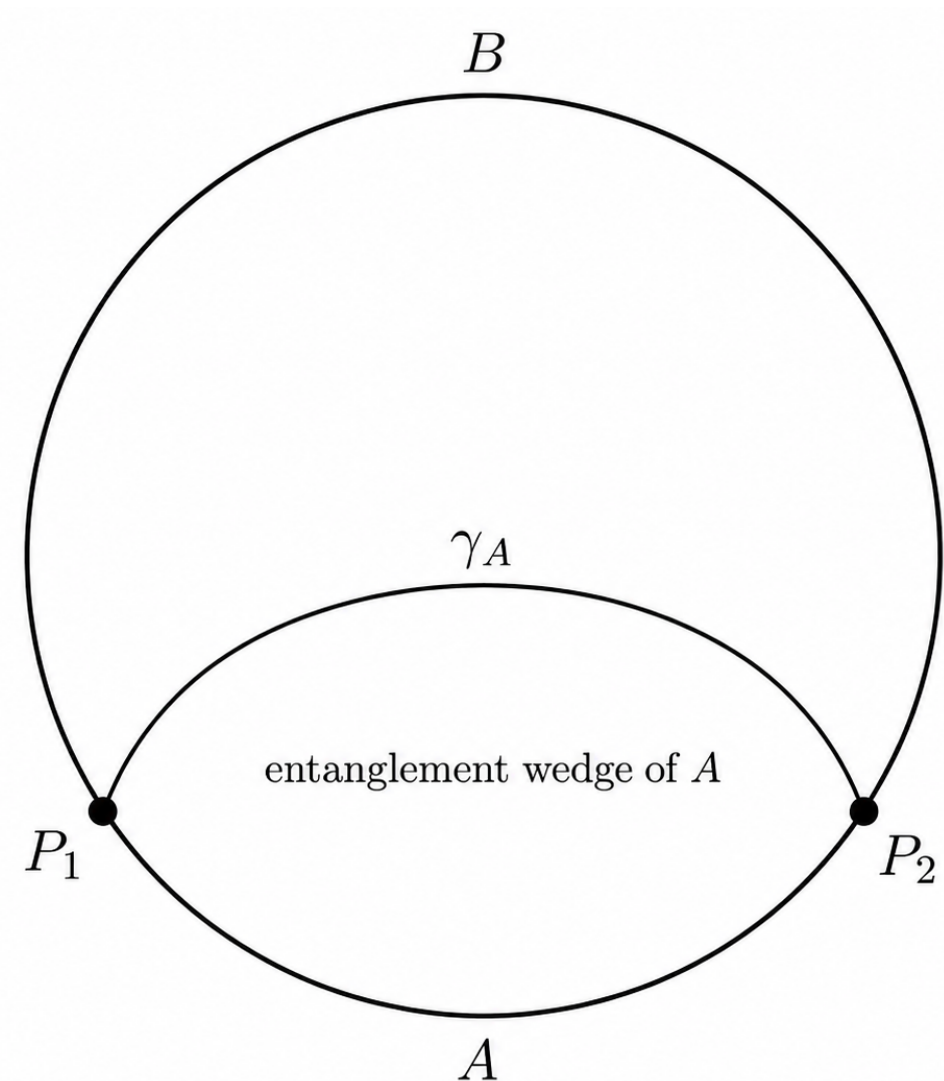
In the Euclidean path integral,  $\rho_A$  is prepared by cutting open region  $A$ . The trace  $\text{Tr} \rho_A^n$  cyclically reconnects these cuts among  $n$  replicas. [1]



**Figure 1:** Path-integral representation of the reduced density matrix and the corresponding  $n$ -sheeted Riemann surface. Cyclic identification of the cuts along  $A$  implements  $\text{Tr} \rho_A^n$ ; here  $n = 3$ .

**Holographic dual.** In holographic CFTs, this boundary replica construction has a bulk geometric saddle. In the  $n \rightarrow 1$  limit, the entropy is given by the Ryu-Takayanagi surface

$$S(A) = \frac{\text{Area}(\gamma_A)}{4G_N}, \quad \partial \gamma_A = \partial A.$$

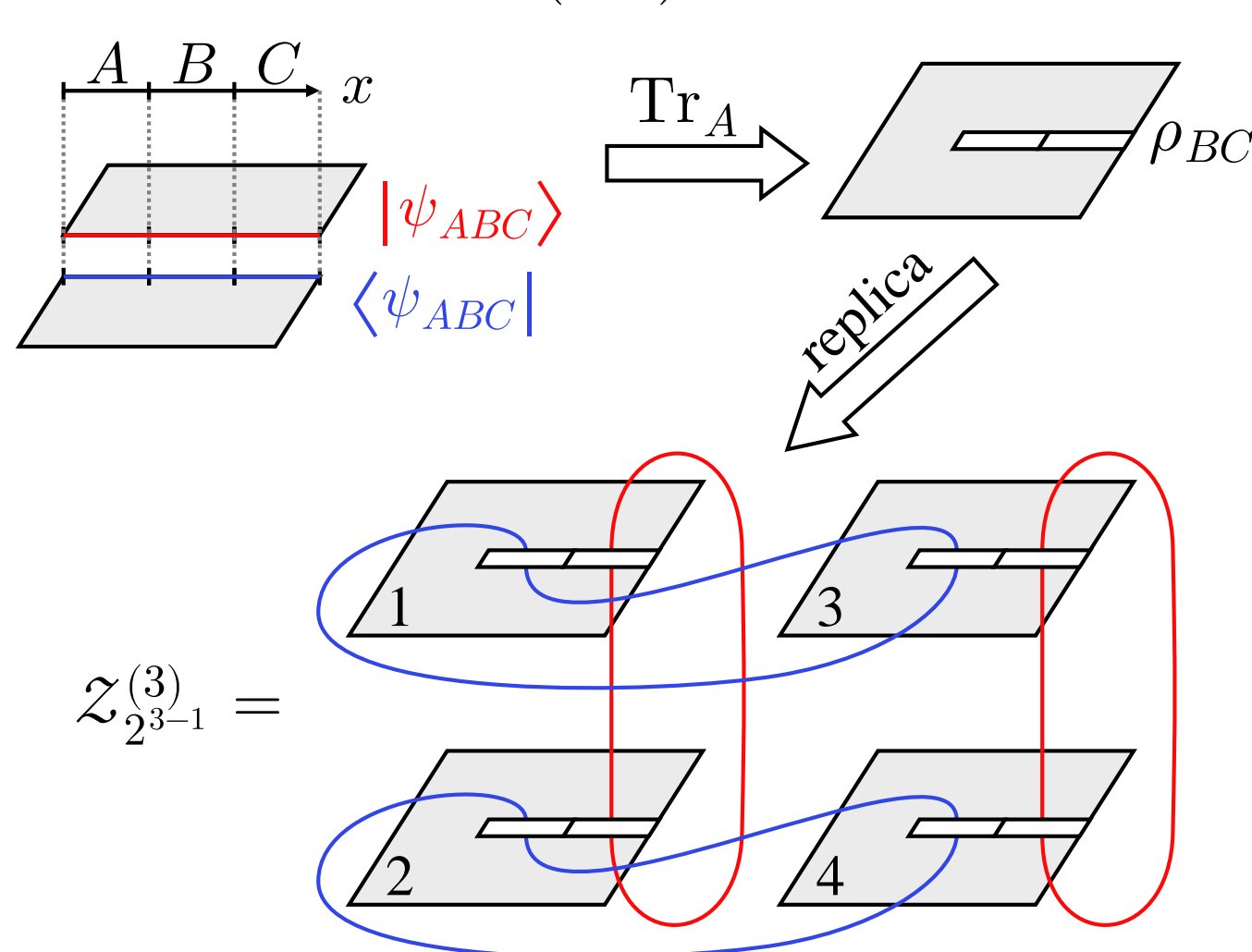


**Figure 2:** Holographic dual of entanglement entropy. The interval  $A$  is anchored at boundary points  $P_1, P_2$ , and its entropy is computed by the bulk geodesic  $\gamma_A$ . The enclosed bulk domain is the entanglement wedge of  $A$ .

## 2. Multientropy and minimal web

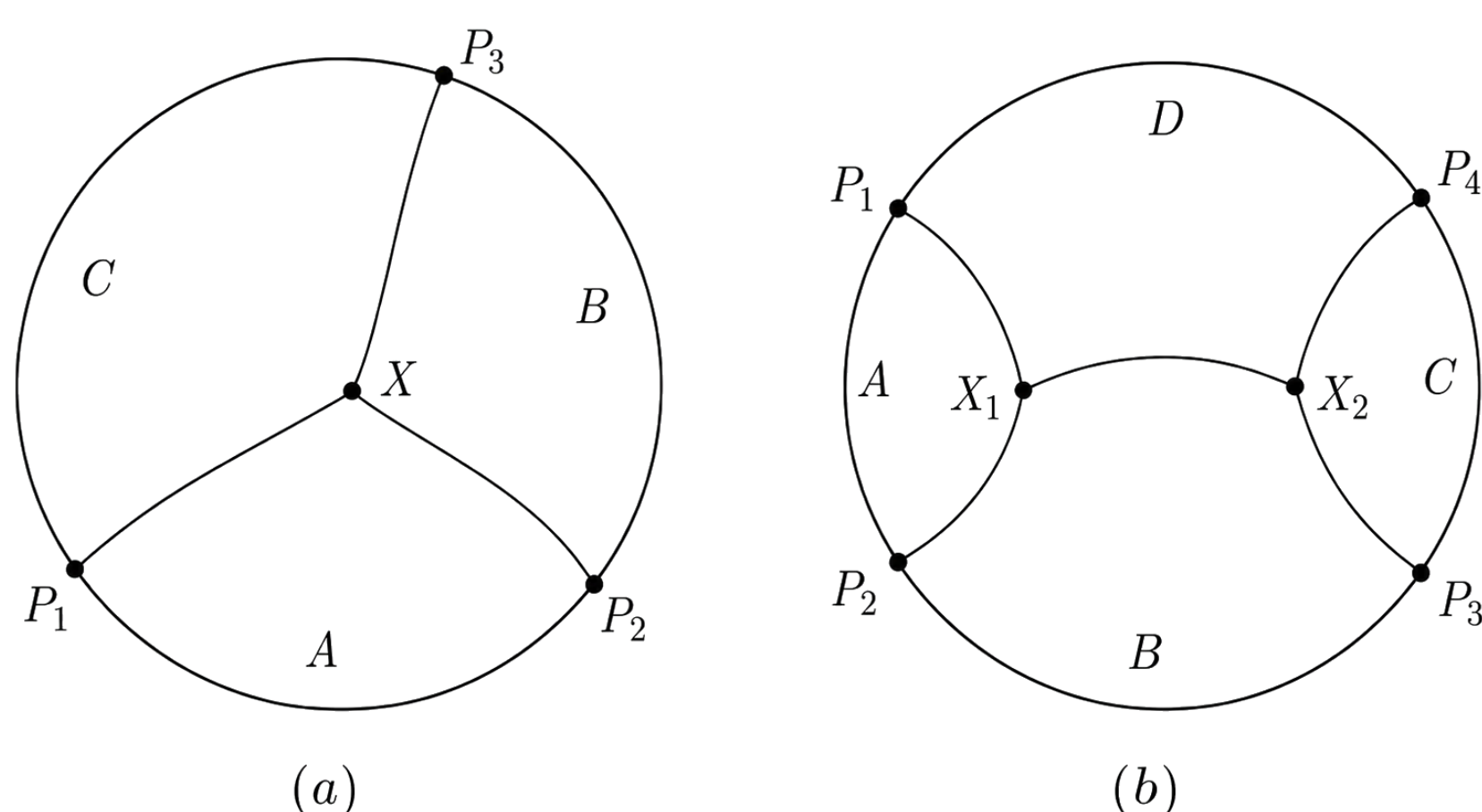
Multientropy generalizes the replica construction by assigning independent permutation patterns to different subsystems. For a  $q$ -partite state,

$$S_n^{(q)} = \frac{1}{(1-n)n^{q-2}} \log \left[ \frac{Z_{n^{q-1}}^{(q)}}{(Z_1^{(q)})^{n^{q-1}}} \right], \quad S^{(q)} = \lim_{n \rightarrow 1} S_n^{(q)}.$$



**Figure 3:** Tripartite multientropy for  $q = 3, n = 2$ . The four replicas form a two-dimensional replica pattern: different colors encode independent subsystem permutations.

For pure holographic states, the RT surface is replaced by a minimal surface web [2].

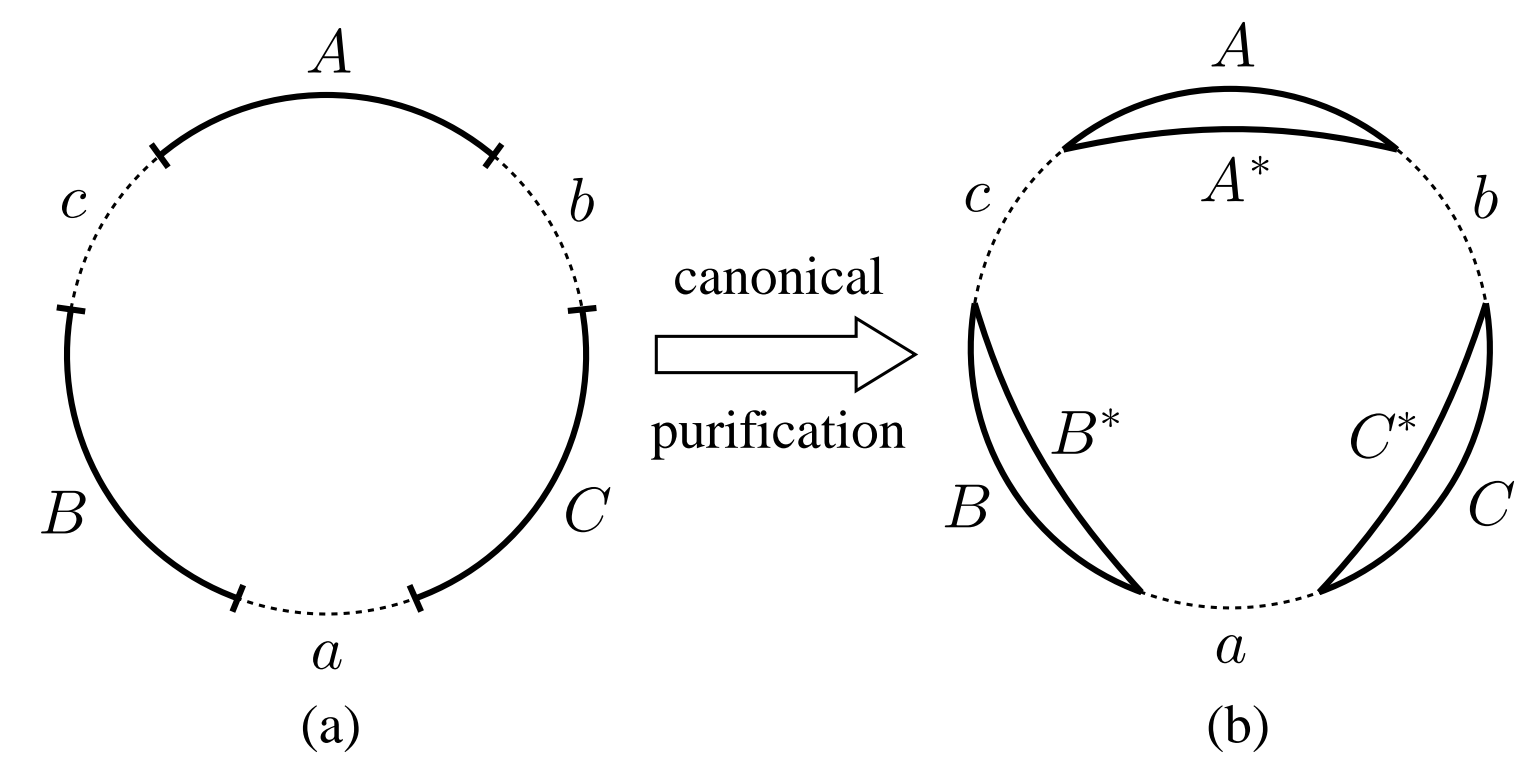


**Figure 4:** Minimal-web dual of multientropy. (a) A tripartite web ending on the boundary points  $P_1, P_2, P_3$ , with boundary regions  $A, B, C$ . (b) A four-partite web with two bulk junctions  $X_1, X_2$ , separating the regions  $A, B, C, D$ .

## 3. Canonical purification

Multientropy is naturally a pure-state construction. To extend it to a mixed state  $\rho_{ABC}$ , we use canonical purification [3]:

$$|\sqrt{\rho_{ABC}}\rangle \in (\mathcal{H}_A \otimes \mathcal{H}_{A^*}) \otimes (\mathcal{H}_B \otimes \mathcal{H}_{B^*}) \otimes (\mathcal{H}_C \otimes \mathcal{H}_{C^*}).$$



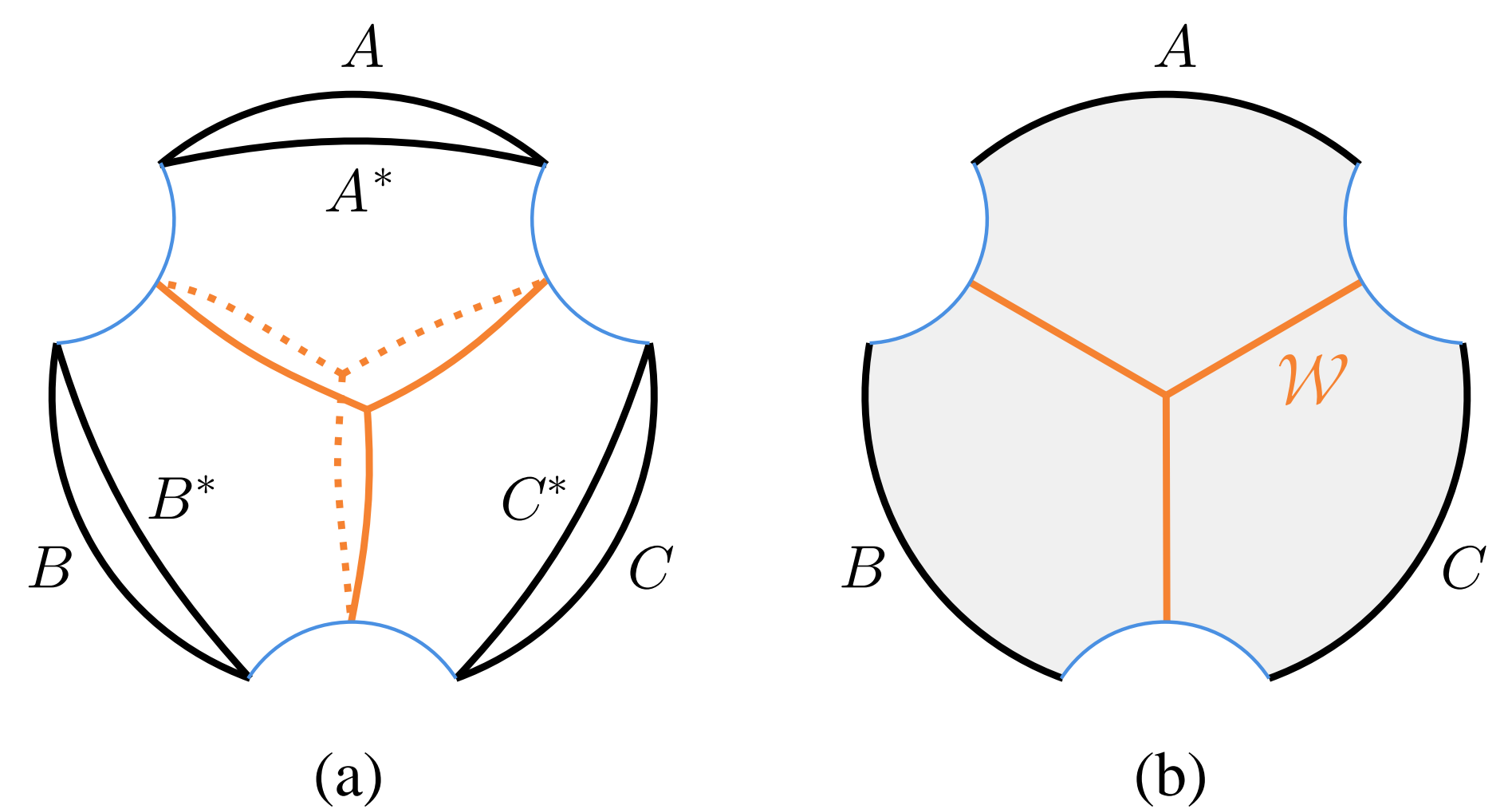
**Figure 5:** Canonical purification of  $\rho_{ABC}$ . A second copy is glued along the traced-out degrees of freedom, producing the purified state  $|\sqrt{\rho_{ABC}}\rangle$ .

## 4. Reflected multientropy and minimal web

We define the reflected multientropy as the multientropy of the canonically purified state [4]:

$$S_R^{(3)}(A:B:C) = S^{(3)}(AA^*:BB^*:CC^*)|_{\sqrt{\rho_{ABC}}}.$$

It reduces to the reflected entropy for  $q = 2$ , and to twice the multientropy for pure states.



**Figure 6:** Holographic canonical purification. The bulk dual is obtained by doubling the entanglement wedge and gluing along the RT surface. In one copy,  $S_R^{(3)}$  is computed by a minimal web  $W$  anchored on the RT surface.

The holographic proposal is

$$S_R^{(q)} = \frac{2}{4G_N} \min_W L[W].$$

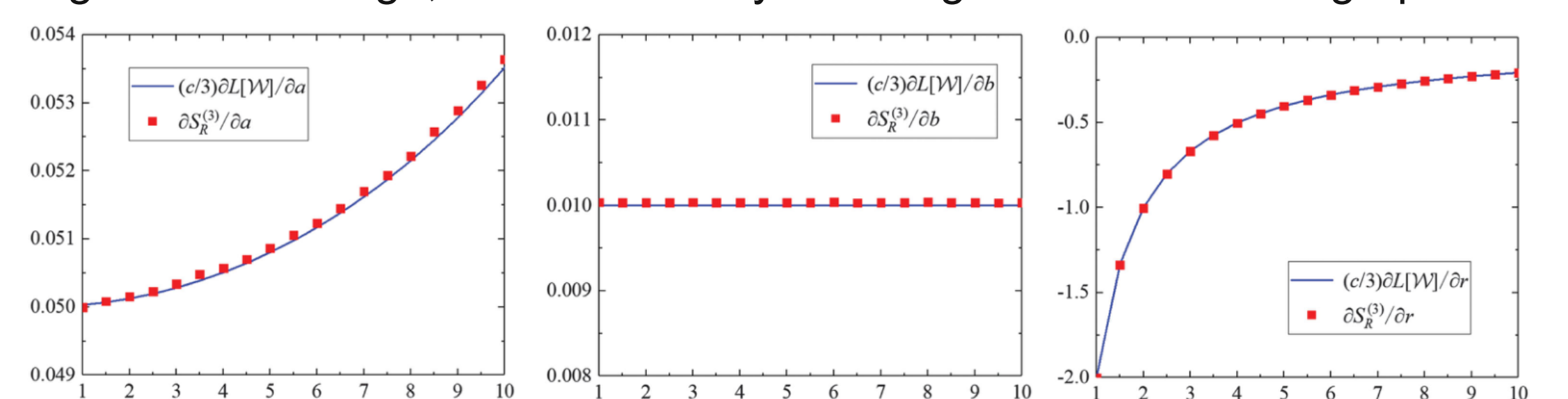
The factor of two comes from the two copies in the canonical purification, while  $L[W]$  is evaluated in a single entanglement wedge.

## 5. Check in AdS<sub>3</sub>/CFT<sub>2</sub>

For the tripartite case, the replicated partition function is a six-point twist-operator correlator,

$$Z_{n^{q-1}, m} = \langle \sigma_1 \sigma_2 \sigma_3 \sigma_4 \sigma_5 \sigma_6 \rangle_{\text{CFT}^{\otimes mn^{q-1}}}.$$

At large central charge, the monodromy result agrees with the holographic web.



**Figure 7:** Zero-temperature check in  $\text{AdS}_3/\text{CFT}_2$ . The derivatives of the CFT result for  $S_R^{(3)}$  match the derivatives of the holographic web length  $(c/3)L[W]$  with high precision.

The same comparison can be extended to finite temperature by putting the CFT on a thermal cylinder and evaluating the bulk web in the BTZ background.

## References

- [1] S. Ryu and T. Takayanagi, "Aspects of holographic entanglement entropy," *JHEP* **08** (2006) 045, arXiv:hep-th/0605073.
- [2] A. Gadde, V. Krishna and T. Sharma, "Towards a classification of holographic multi-partite entanglement measures," *JHEP* **08** (2023) 202, arXiv:2304.06082, doi:10.1007/JHEP08(2023)202.
- [3] S. Dutta and T. Faulkner, "A canonical purification for the entanglement wedge cross-section," *JHEP* **03** (2021) 178, arXiv:1905.00577, doi:10.1007/JHEP03(2021)178.
- [4] M.-K. Yuan, M. Li and Y. Zhou, "Reflected Multientropy and Its Holographic Dual," *Phys. Rev. Lett.* **135** (2025) 091604, doi:10.1103/PhysRevLett.135.091604.