



Electromagnetic response of disordered superconductors with finite-momentum pairing

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Overview

Background: Finite-momentum pairing with momentum \vec{q} enables photon absorption and optical activity of collective modes.

Questions: 1) How can the electromagnetic response be calculated in a conserving and gauge-invariant way?

2) How does disorder reshape quasiparticle and collective-mode contributions?

Method: Baym-Kadanoff conserving approximation

Results: All BCS-level fluctuations can make non-negligible contributions, demonstrating the necessity of vertex corrections.

Qualitative Analysis

Symmetry analysis

• Quasiparticle part:

$\vec{q} = \vec{0}$: electrons satisfy $\vec{v}_{\vec{k}} = -\vec{v}_{-\vec{k}}$, no net current

finite \vec{q} : $\vec{v}_{\vec{k}+\vec{q}} \neq -\vec{v}_{-\vec{k}+\vec{q}} \Rightarrow$ finite current

• Collective modes part:

s-wave pairing case: (Similar for general pairing symmetry)

$\vec{q} = \vec{0}$: Δ is a scalar, its fluctuation couples via $\vec{\nabla} \cdot \vec{A}$

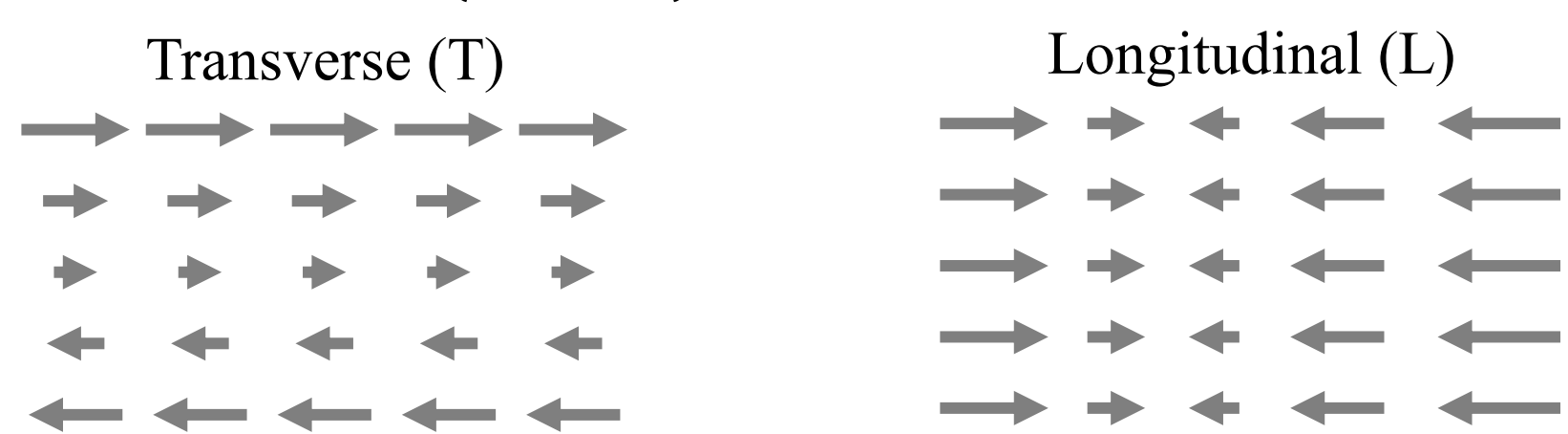
\Rightarrow no contribution in transverse gauge ($\vec{\nabla} \cdot \vec{A} = 0$)

finite \vec{q} : extra scalar $\vec{q} \cdot \vec{A}$ enables finite current

Magnetization analogy

• Quasiparticle current $\Rightarrow \vec{M}$, backflow $\Rightarrow \vec{H}$

net current $\Rightarrow \vec{B} \propto (\vec{M} + \vec{H})$



$$\vec{\nabla} \cdot \vec{H} = 0, \vec{\nabla} \times \vec{H} = 0 \Rightarrow \vec{H} = \vec{0} \quad \vec{\nabla} \cdot \vec{B} = 0, \hat{k} \parallel \vec{M} \parallel \vec{H} \Rightarrow \vec{B} = \vec{0}$$

• $\vec{H} = \vec{0}$: no contribution from collective modes

$\vec{B} = \vec{0}$: exact cancellation of quasiparticle and collective modes

	Galilean invariance	Broken Galilean invariance
$\vec{q} = \vec{0}$	$\vec{M} = \vec{0} \Rightarrow \vec{H} = \vec{0} \Rightarrow \vec{B} = \vec{0}$	
$\vec{q} \neq \vec{0}$	$\vec{M} \neq \vec{0}$ and pure L $\Rightarrow \vec{B} = \vec{0}$	$\vec{M} \neq \vec{0}$ but mixed T and L $\Rightarrow \vec{H} = -\vec{M}_L \Rightarrow \vec{B} \propto \vec{M}_T$

Finite \vec{q} enables optical activity of collective mode.
Broken Galilean invariance prevents exact cancellation.

Framework

Baym-Kadanoff approximation

(with self-consistent Hartree-Fock and Born approximations)

H_{int} : effective e-e attractive interaction + static nonmagnetic disorder potential

Dyson equation

$$G^{-1} = G_0^{-1} + \Sigma$$

Luttinger-Ward functional

$$\Phi[G, \{v\}] = \text{diagrams}$$

Φ -derivable self-energy

$$\Sigma = \text{diagrams}$$

Two-particle irreducible vertex

$$K = \text{diagrams}$$

Bethe-Salpeter equation

$$\Gamma = \gamma + \Gamma K$$

Linear electromagnetic response

$$\text{diagrams}$$

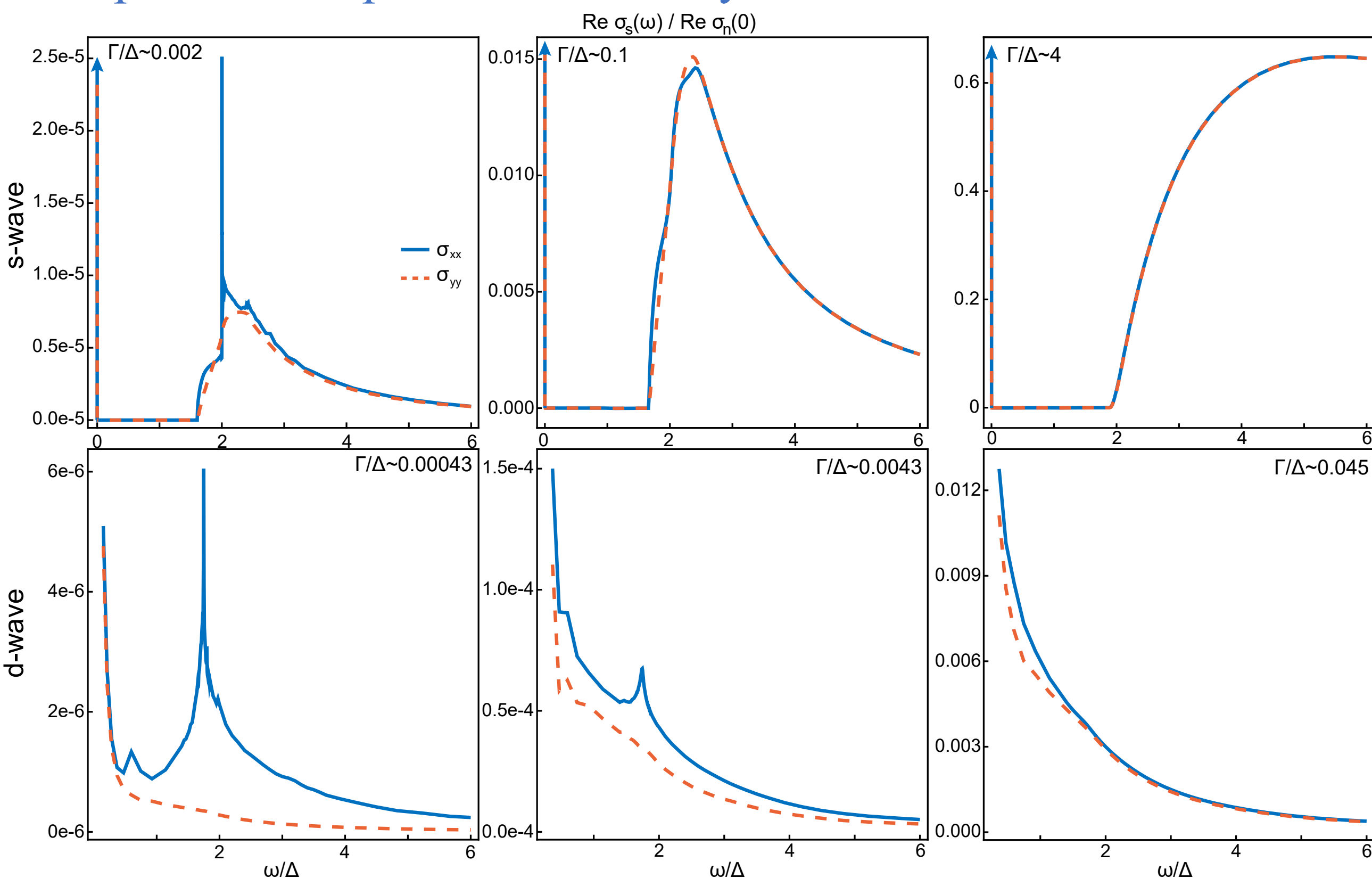
• Equivalent to path-integral approach and gauge-invariant kinetic theory.

• Conserving method including all BCS-level fluctuations.

• More complete than consistent-fluctuation of order parameter method.

Results

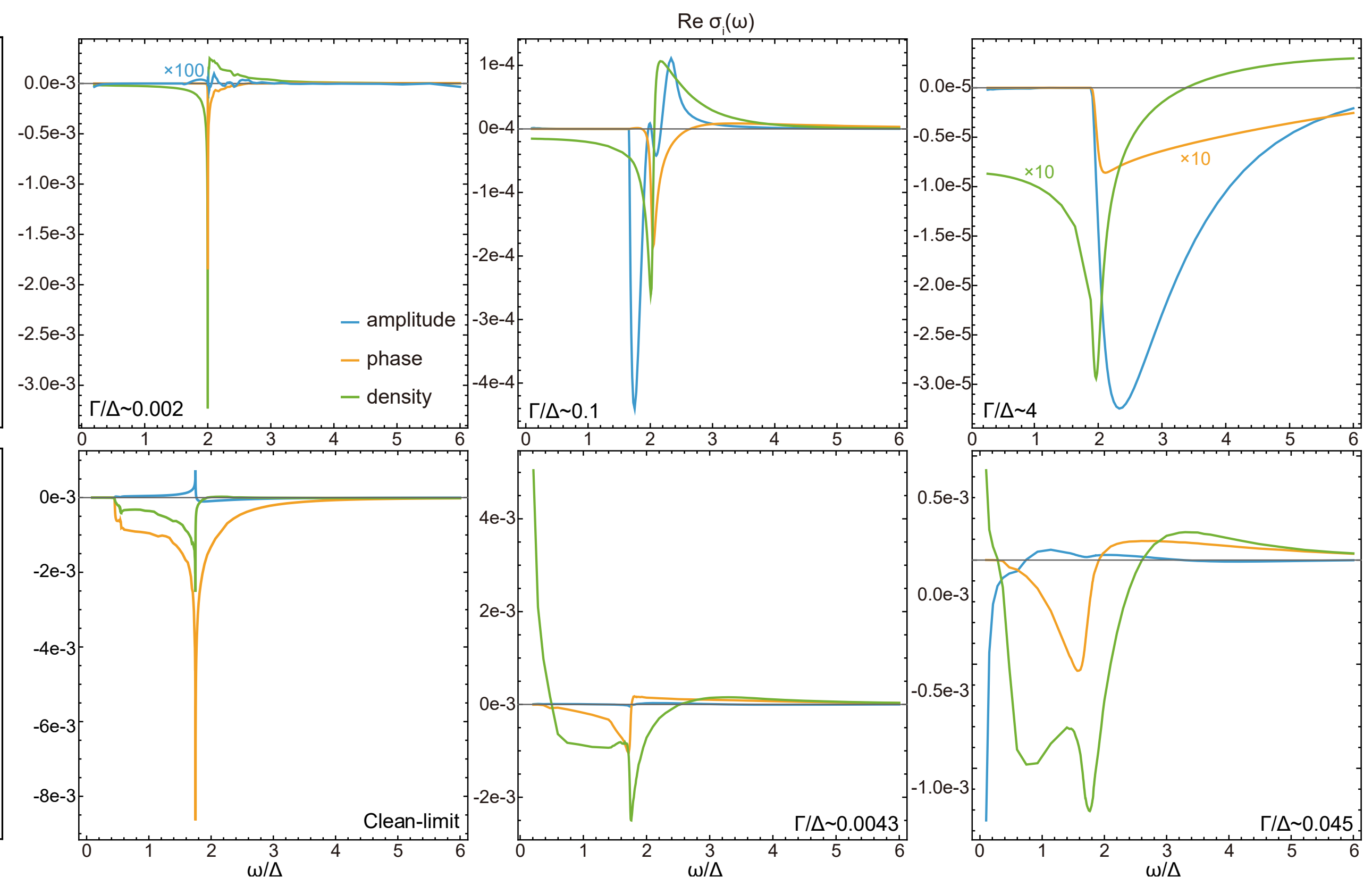
Real part of the optical conductivity ($\vec{q} \parallel \hat{x}$)



• Cleaner case: peaks reflect quasiparticle DOS; strongly anisotropic

• Dirtier case: peaks suppressed; more isotropic; approach the Mattis-Bardeen limit

Contribution from different collective modes



• Amplitude/density contributions grow relative to phase contribution with disorder

• Vertex corrections are important near optical-response peaks

References

[1] PRB 108, 224516 (2023)

[2] arXiv:2410.1867 & arXiv:2501.13722

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