

## Introduction

### ► Harmonic Voltage Measurement

Under a current excitation  $I$ , the resistance may depend on the current

$$R(I) = R_0 + R_1 I + R_2 I^2 + \dots$$

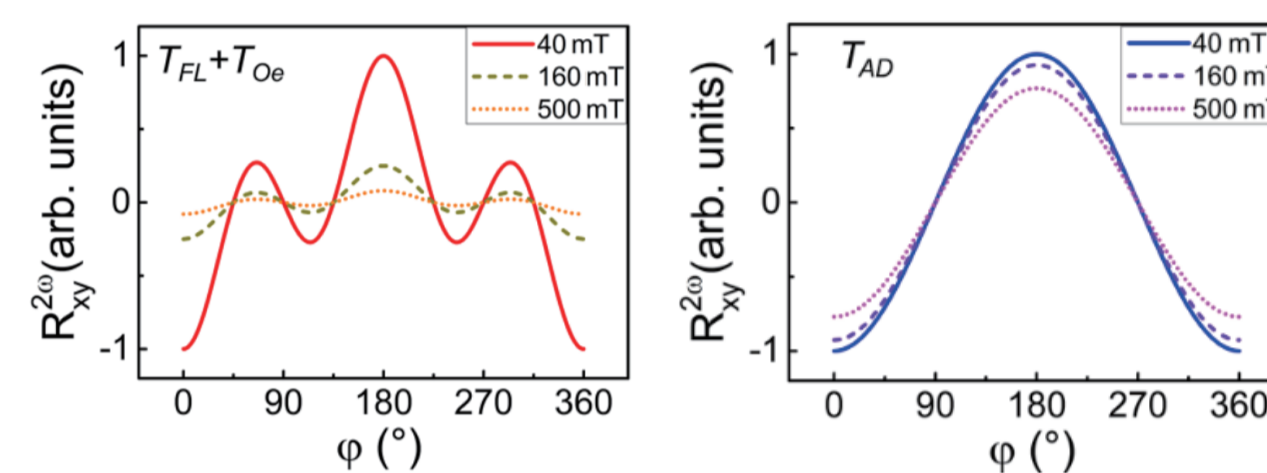
For sinusoidal current  $I = I_0 \sin(\omega t)$ ,

$$V = IR(I) = R_0 I_0 \sin(\omega t) + R_1 I_0^2 \sin^2(\omega t) + R_2 I_0^3 \sin^3(\omega t) + \dots = V_{DC} + V_{1\omega} \sin(\omega t) + V_{2\omega} \cos(2\omega t) + V_{3\omega} \sin(3\omega t) + \dots$$

$V_{DC}$ : rectified voltage  
 $V_{1\omega}$ : first harmonic voltage  
 $V_{2\omega}$ : second harmonic voltage  
 $V_{3\omega}$ : third harmonic voltage  
 ...

### ► Second Harmonic Measurement

- Spin-orbital torque (SOT)
- Unidirectional spin Hall magnetoresistance (USMR)
- Thermal effects (anomalous Nernst effects, spin Seebeck effect...)



$$R_{xy}^2\omega(\phi) = \{ [R_{AHE}^{(B_{DLT}/B_{eff})} + R_{VIT} \cos\phi + 2R_{PHE}^{(B_{DLT} + B_{Oe})/B_{eff}} ] 2\cos^3\phi - \cos\phi \}$$

Can Onur Avci et al, Physical Review B 90, 224427 (2014)

### ► Third Harmonic Measurement

- SOT: nonlinear term of magnetization oscillation
- SOT: nonlinear term of Spin Hall effect (NSHE)
- Quantum geometry effects
- Magnetoelastic effect
- Thermally induced magnetization reduction

**A new mechanism: thermally induced magnetocrystalline anisotropy variation**

## Thermal Effects in 3rd Harmonic Measurement: derivation

### ► Derivation

$$R = R(B, T_0) + \frac{\partial R(B, T_0)}{\partial T} \delta T = R(B, T_0) + \alpha \frac{\partial R(B, T_0)}{\partial T} I^2$$

for  $I = I_0 \sin(\omega t)$ :

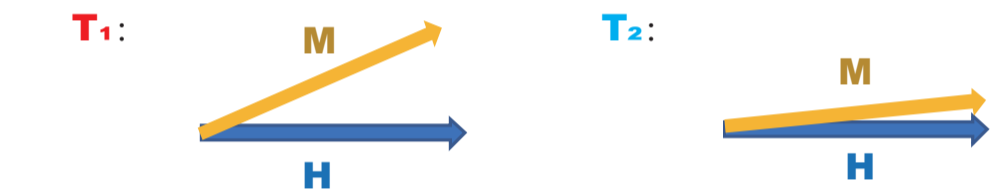
$$V = IR = I_0 R \sin(\omega t) + \alpha \frac{\partial R(B, T_0)}{\partial T} I_0^3 \sin^3(\omega t) = \left( I_0 R + \frac{3}{4} \alpha \frac{\partial R(B, T_0)}{\partial T} I_0^3 \right) \sin(\omega t) - \alpha \frac{1}{4} \frac{\partial R(B, T_0)}{\partial T} I_0^3 \sin(3\omega t)$$

for  $R_{xx} = R_0 + R_{AMR} \cos 2\phi_M$ ,  $R_{xy} = R_{PHE} \sin 2\phi_M$

$$\frac{\partial R_{xx}}{\partial T} = \frac{\partial R_0}{\partial T} + \frac{\partial R_{AMR}}{\partial T} \cos 2\phi_M - 2R_{AMR} \sin 2\phi_M \frac{\partial \phi_M}{\partial T}$$

$$\frac{\partial R_{xy}}{\partial T} = \frac{\partial R_{PHE}}{\partial T} \sin 2\phi_M + 2R_{PHE} \cos 2\phi_M \frac{\partial \phi_M}{\partial T}$$

The derivative  $\partial\phi/\partial T$  can exhibit rich angular dependencies when influenced by magnetocrystalline anisotropy



for in-plane uniaxial magnetocrystalline anisotropy:  $E = M \cdot H - H_{Ku} M_s / 2 \sin^2 \phi_M$

$$\frac{\partial \phi_M}{\partial T} = \frac{\partial H_{Ku}}{\partial T} \left( \frac{\sin 2\phi_M}{2H} + \frac{H_{Ku}}{2H^2} \sin 2\phi_M \cos 2\phi_M \right)$$

for in-plane 4-fold magnetocrystalline anisotropy:  $E = M \cdot H - H_{Kc} M_s / 2 \sin^2 \phi_M \cos^2 \phi_M$

$$\frac{\partial \phi_M}{\partial T} = \frac{\partial H_{Kc}}{\partial T} \left( \frac{\sin 4\phi_M}{4H} + \frac{H_{Kc}}{4H^2} \sin 4\phi_M \cos 4\phi_M \right)$$

### ► Results

#### Uniaxial Magnetocrystalline Anisotropy

$$R_{xx}^{3\omega} = R_0^{xx,3\omega} + R_{AMR}^{xx,3\omega} \cos 2\phi_M - R_{\Delta H_{Ku}1}^{xx,3\omega} \sin^2 2\phi_M - R_{\Delta H_{Ku}2}^{xx,3\omega} \sin^2 2\phi_M \cos 2\phi_M$$

$$R_{xy}^{3\omega} = R_{AMR}^{xy,3\omega} \sin 2\phi_M + R_{\Delta H_{Ku}1}^{xy,3\omega} \sin 2\phi_M \cos 2\phi_M + R_{\Delta H_{Ku}2}^{xy,3\omega} \sin 2\phi_M \cos 2\phi_M \cos 2\phi_M$$

where

$$R_0^{xx,3\omega} = \frac{1}{4} \alpha I_0^2 \frac{\partial R_0}{\partial T}$$

$$R_{AMR}^{xx,3\omega} = \frac{1}{4} \alpha I_0^2 \frac{\partial R_{AMR}}{\partial T}$$

$$R_{\Delta H_{Ku}1}^{xx,3\omega} = \frac{1}{4} \alpha I_0^2 \cdot R_{AMR} \frac{\partial H_{Ku}}{\partial T} \frac{1}{H}$$

$$R_{\Delta H_{Ku}2}^{xx,3\omega} = \frac{1}{4} \alpha I_0^2 \cdot R_{AMR} \frac{\partial H_{Ku}}{\partial T} \frac{H_{Ku}}{H^2}$$

$$R_{\Delta H_{Ku}1}^{xy,3\omega} = \frac{1}{4} \alpha I_0^2 \cdot R_{PHE} \frac{\partial H_{Ku}}{\partial T} \frac{1}{H}$$

$$R_{\Delta H_{Ku}2}^{xy,3\omega} = \frac{1}{4} \alpha I_0^2 \cdot R_{PHE} \frac{\partial H_{Ku}}{\partial T} \frac{H_{Ku}}{H^2}$$

#### 4-fold Magnetocrystalline Anisotropy

$$R_{xx}^{3\omega} = R_0^{xx,3\omega} + R_{AMR}^{xx,3\omega} \cos 2\phi_M - R_{\Delta H_{Kc}1}^{xx,3\omega} \sin 2\phi_M \sin 4\phi_M - R_{\Delta H_{Kc}2}^{xx,3\omega} \sin 2\phi_M \sin 4\phi_M \cos 4\phi_M$$

$$R_{xy}^{3\omega} = R_{AMR}^{xy,3\omega} \sin 2\phi_M + R_{\Delta H_{Kc}1}^{xy,3\omega} \cos 2\phi_M \sin 4\phi_M + R_{\Delta H_{Kc}2}^{xy,3\omega} \cos 2\phi_M \sin 4\phi_M \cos 4\phi_M$$

where

$$R_0^{xx,3\omega} = \frac{1}{4} \alpha I_0^2 \frac{\partial R_0}{\partial T}$$

$$R_{AMR}^{xx,3\omega} = \frac{1}{4} \alpha I_0^2 \frac{\partial R_{AMR}}{\partial T}$$

$$R_{\Delta H_{Kc}1}^{xx,3\omega} = \frac{1}{4} \alpha I_0^2 \cdot R_{AMR} \frac{\partial H_{Kc}}{\partial T} \frac{1}{2H}$$

$$R_{\Delta H_{Kc}2}^{xx,3\omega} = \frac{1}{4} \alpha I_0^2 \cdot R_{AMR} \frac{\partial H_{Kc}}{\partial T} \frac{H_{Kc}}{2H^2}$$

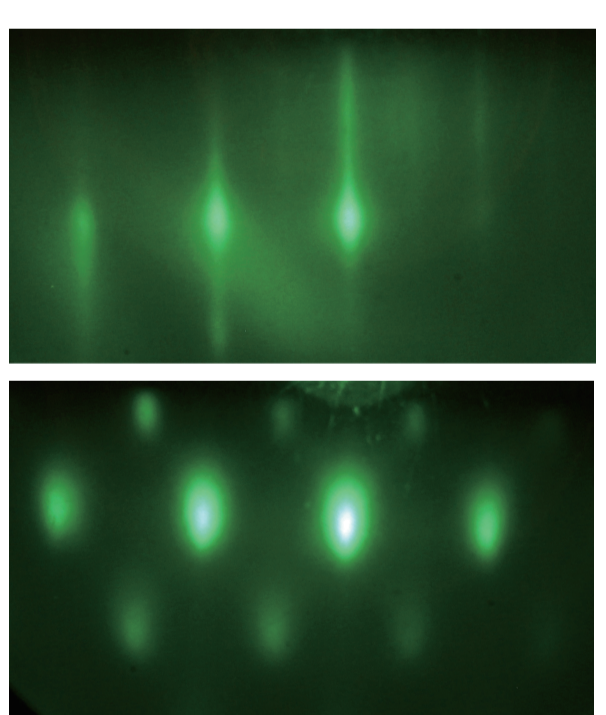
$$R_{\Delta H_{Kc}1}^{xy,3\omega} = \frac{1}{4} \alpha I_0^2 \cdot R_{PHE} \frac{\partial H_{Kc}}{\partial T} \frac{1}{2H}$$

$$R_{\Delta H_{Kc}2}^{xy,3\omega} = \frac{1}{4} \alpha I_0^2 \cdot R_{PHE} \frac{\partial H_{Kc}}{\partial T} \frac{H_{Kc}}{2H^2}$$

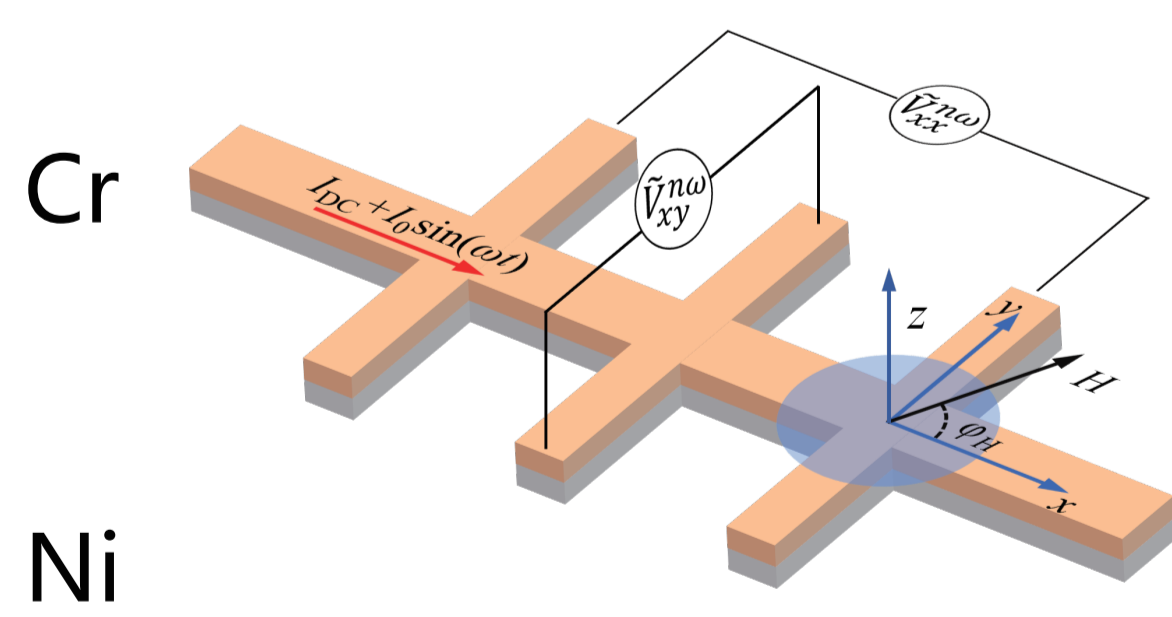
## 3rd Harmonic signal in Cr/Ni: 4 Fold Anisotropy

### MgO(001)/Cr(7nm)/Ni(7nm)

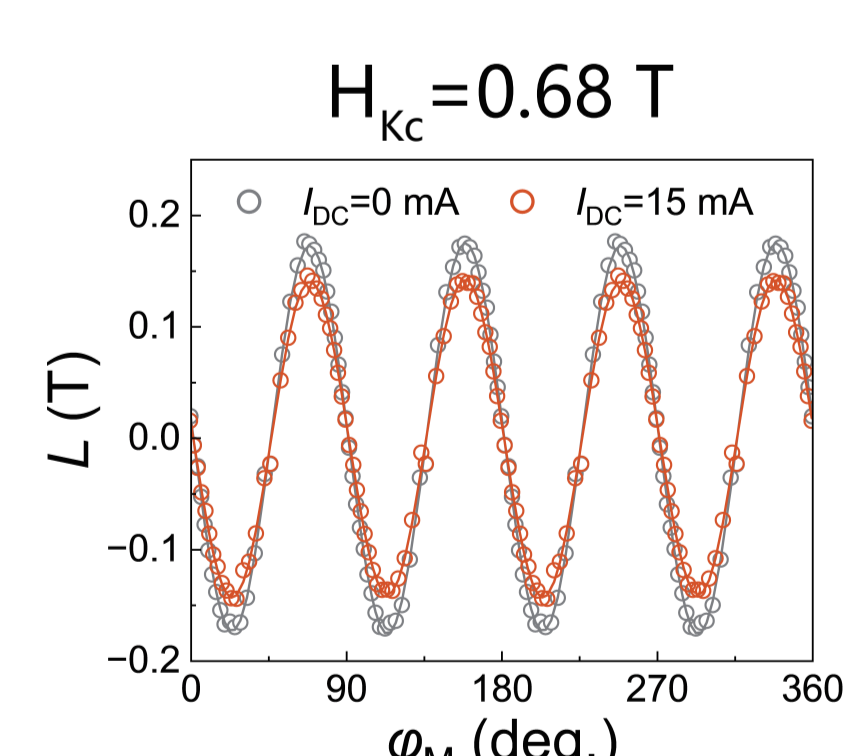
RHEED pattern



Hall bar measurement

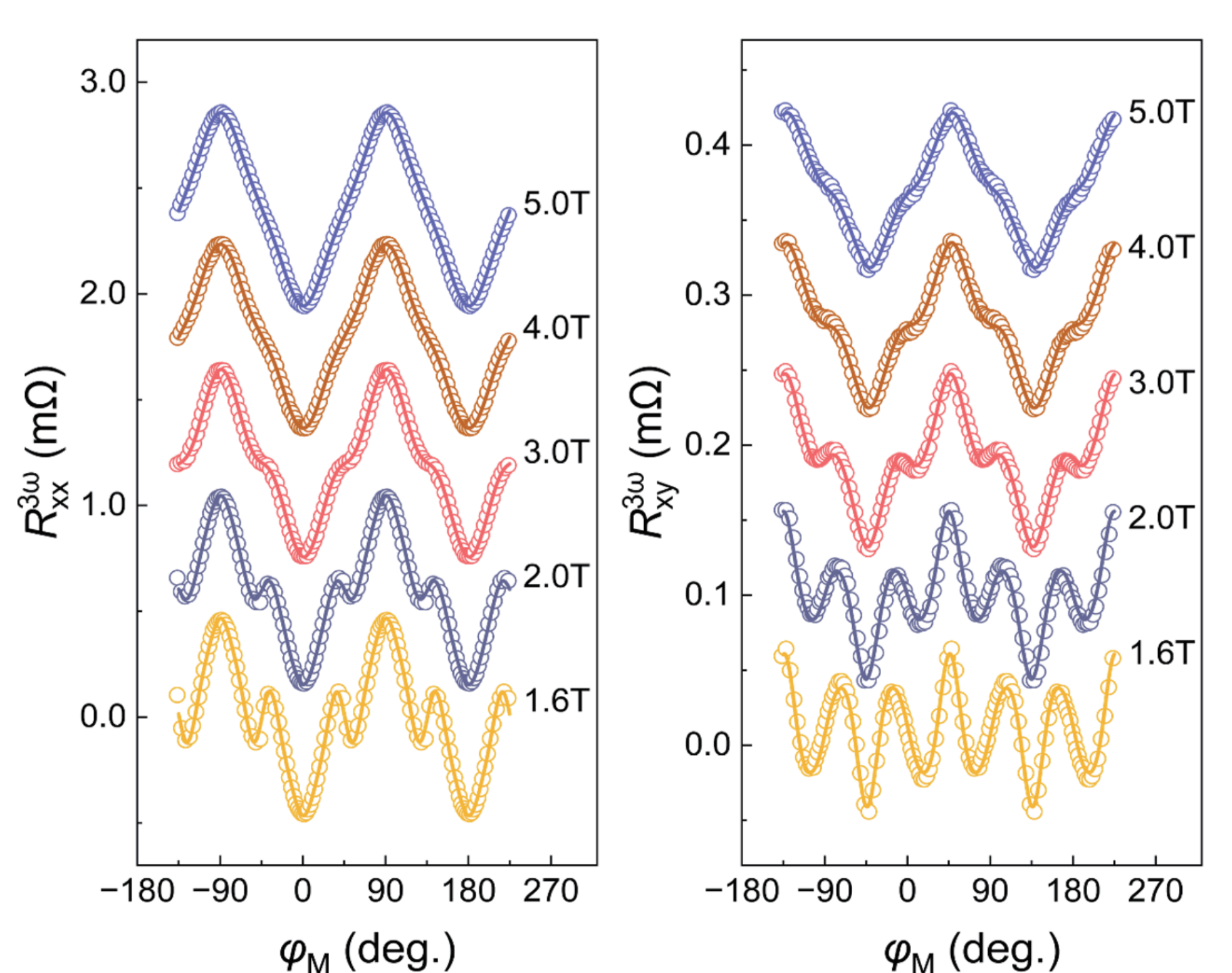


$I$  is along the Cr(001)[100]

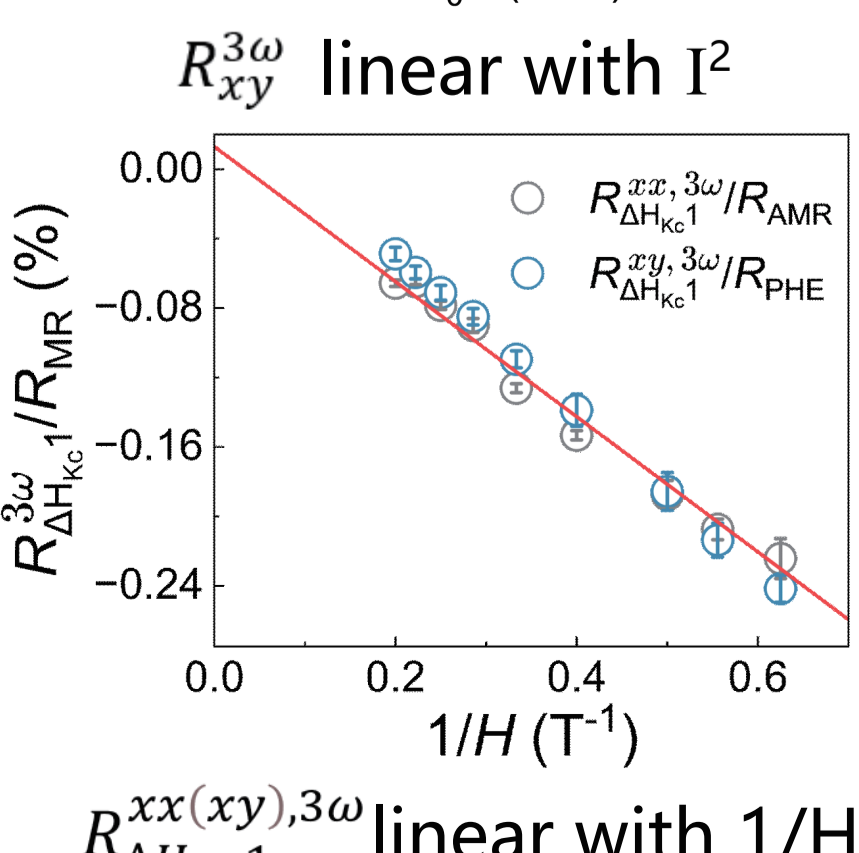
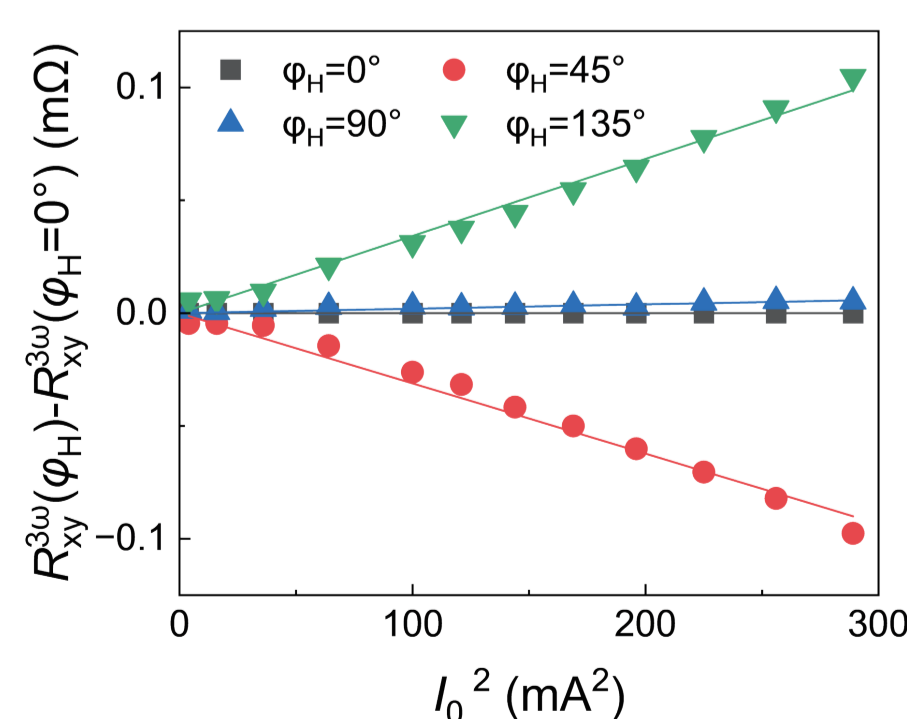


strong 4 fold magnetocrystalline anisotropy

### ► 3rd Harmonic Resistance ( $I = 15$ mA)



strong  $\cos(6\phi)$  or  $\sin(6\phi)$  signal, suppressed in high H

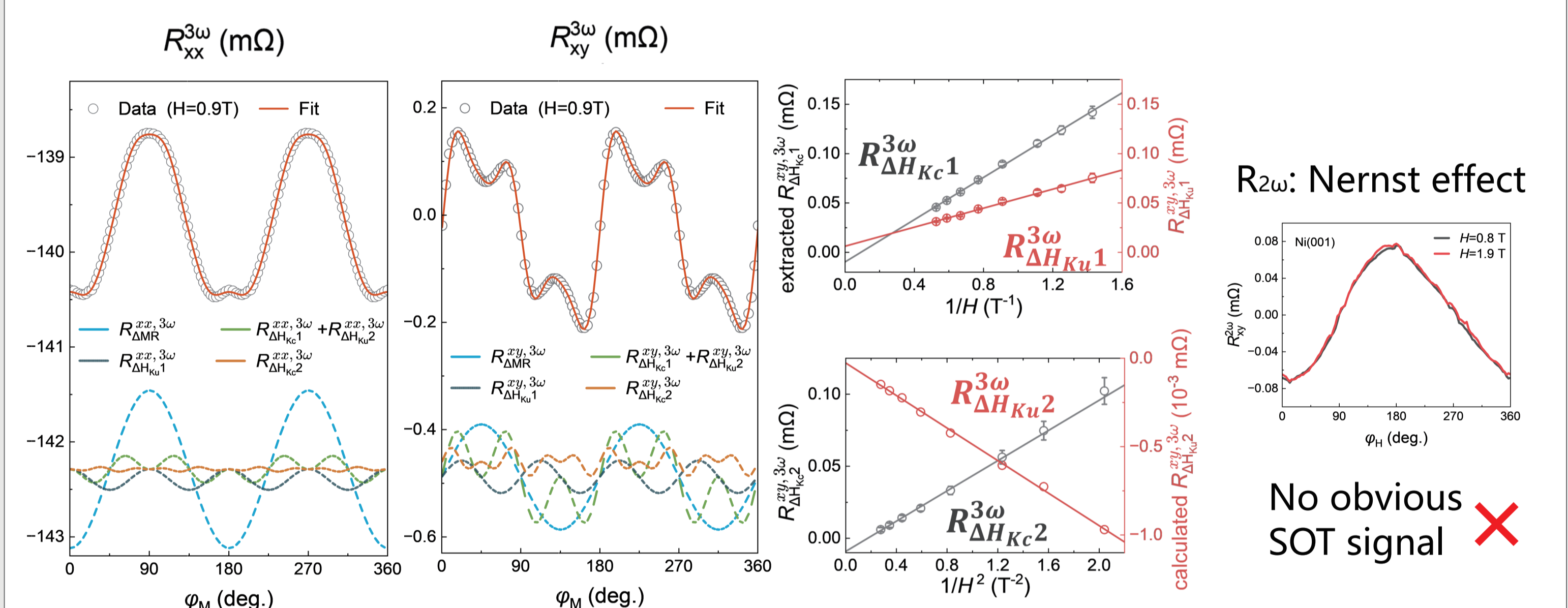


## 3rd Harmonic in Ni

### ► Single Crystal Ni MgO(001)/Ni(10nm)

uniaxial anisotropy + 4 fold anisotropy ✓

$R_{3\omega}$ ,  $I = 12$  mA,  $H = 0.7$  T



both  $R_{\Delta H_{Kc}}^{3\omega}$  and  $R_{\Delta H_{Ku}}^{3\omega}$  dominate different field scaling

### Anisotropy measurement

$I = I_{DC} + I_0 \sin(\omega t)$ ,  $I_0 = 2$  mA

$H_{ku}$ : Uniaxial Anisotropy Field  $H_{kc}$ : 4-fold Anisotropy Field

#### ► change offset current

$H_{ku} \propto I_{DC}^2$ , independent of sample temperature

$H_{kc} \propto I_{DC}^2$  and temperature

#### ► change temperature

$H_{ku}$  (kOe) vs  $I_{DC}^2$  (mA<sup>2</sup>)

$H_{kc}$  (kOe) vs  $T$  (K)

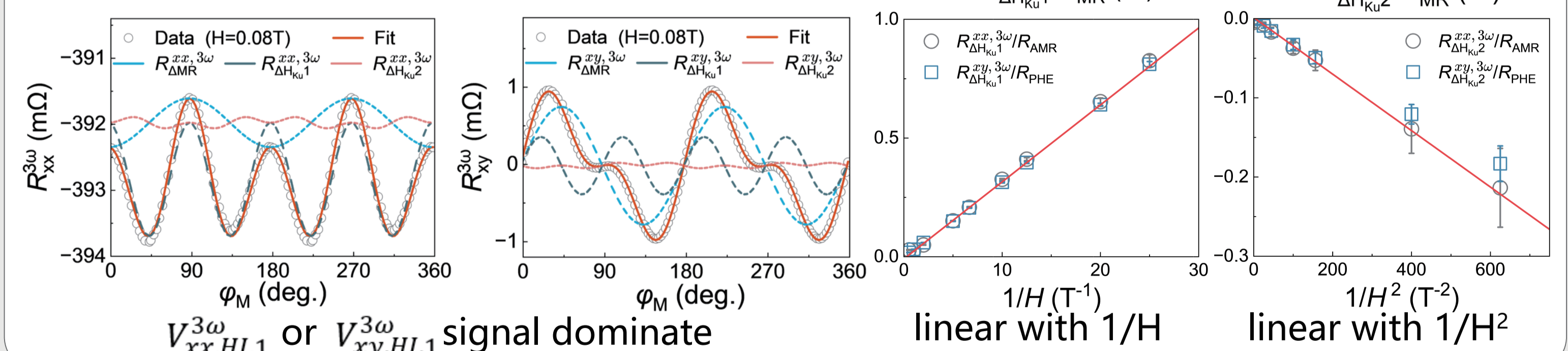
★  $H_{ku} \propto I_{DC}^2$ , independent of sample temperature

$H_{kc} \propto I_{DC}^2$  and temperature

### ► Poly Ni Al2O3(0001)/SiO2(4nm)/Ni(10nm)

uniaxial anisotropy ✓

$R_{3\omega}$ ,  $I = 12$  mA,  $H = 8000$  Oe



$V_{xx,HI,1}^{3\omega}$  or  $V_{xy,HI,1}^{3\omega}$  signal dominate

linear with  $1/H$

linear with  $1/H^2$

## Summary

- Third-harmonic voltage in Cr/Ni and Ni systems
- New mechanism identified: thermally induced magnetocrystalline anisotropy variation
- Third-harmonic voltage depends on the type of magnetocrystalline anisotropy in this mechanism
- Thermal effects must be carefully considered in third-harmonic measurements